ilmedia

Philipp, Friedrich

Relatively bounded perturbations of J-non-negative operators

This work is licensed under a [Creative Commons Attribution 4.0](https://creativecommons.org/licenses/by/4.0/) [International](https://creativecommons.org/licenses/by/4.0/) license. To view a copy of this license, visit <https://creativecommons.org/licenses/by/4.0/>

TU Ilmenau | Universitätsbibliothek | ilmedia, 2024 <http://www.tu-ilmenau.de/ilmedia>

Relatively Bounded Perturbations of J-Non-Negative Operators

Friedrich Philipp[1](http://orcid.org/0000-0002-4670-8894)

Received: 29 June 2021 / Accepted: 6 July 2022 / Published online: 21 December 2022 © The Author(s) 2022

Abstract

We improve known perturbation results for self-adjoint operators in Hilbert spaces and prove spectral enclosures for diagonally dominant *J* -self-adjoint operator matrices. These are used in the proof of the central result, a perturbation theorem for *J* -nonnegative operators. The results are applied to singular indefinite Sturm-Liouville operators with L^p -potentials. Known bounds on the non-real eigenvalues of such operators are improved.

Keywords *J* -self-adjoint operator · *J* -non-negative operator · Relatively bounded perturbation · Spectral enclosure

Mathematics Subject Classification Primary 47B50 · 47A55; Secondary 47E05 · 34L15

1 Introduction

J -self-adjoint operators are special generally non-self-adjoint operators in Hilbert spaces; they appear in various applications in mathematics and mathematical physics such as the Klein-Gordon equation [\[17,](#page-29-0) [18](#page-29-1), [24](#page-29-2), [40,](#page-30-0) [41\]](#page-30-1) and other types of wave equations $[14, 21, 30, 31]$ $[14, 21, 30, 31]$ $[14, 21, 30, 31]$ $[14, 21, 30, 31]$ $[14, 21, 30, 31]$ $[14, 21, 30, 31]$ $[14, 21, 30, 31]$, PT-symmetry in quantum mechanics $[8, 32, 42, 43, 46]$ $[8, 32, 42, 43, 46]$ $[8, 32, 42, 43, 46]$ $[8, 32, 42, 43, 46]$ $[8, 32, 42, 43, 46]$ $[8, 32, 42, 43, 46]$ $[8, 32, 42, 43, 46]$ $[8, 32, 42, 43, 46]$ $[8, 32, 42, 43, 46]$ $[8, 32, 42, 43, 46]$ & [\[3](#page-29-6), Sec. 7], and self-adjoint analytic operator functions [\[35](#page-30-8), [37](#page-30-9)[–39\]](#page-30-10), just to name a few.

They also occur naturally as realizations of indefinite Sturm-Liouville expressions [\[5](#page-29-7), [6](#page-29-8), [12,](#page-29-9) [13](#page-29-10), [29,](#page-30-11) [45](#page-30-12), [48](#page-30-13)]. As a motivation, let us consider such an indefinite Sturm-

Communicated by Jussi Behrndt.

 \boxtimes Friedrich Philipp friedrich.philipp@tu-ilmenau.de http://www.tu-ilmenau.de/obc/team/friedrich-philipp

¹ Technische Universität Ilmenau, Institute of Mathematics, Ilmenau, Germany

$$
A: f \mapsto \text{sgn} \cdot (-f'' + qf) \tag{1.1}
$$

with a real-valued potential $q \in L^p(\mathbb{R})$, $p > 1$, defined on the maximal domain

dom
$$
A = \{ f \in L^2(\mathbb{R}) : f, f' \in AC_{loc}(\mathbb{R}), -f'' + qf \in L^2(\mathbb{R}) \}.
$$

It is well known that the non-real spectrum of *A* is bounded and consists of isolated eigenvalues, see, e.g., $[5, 7]$ $[5, 7]$ $[5, 7]$, and it is of particular interest to find bounds on these eigenvalues.

The first result in this direction has been established in [\[6\]](#page-29-8) for $q \in L^{\infty}(\mathbb{R})$. Later, it was refined and immensely generalized to unbounded potentials and a large class of weights in [\[7](#page-29-11)]. Recently, the bound on the imaginary part from [\[7](#page-29-11)] could be further improved in [\[9](#page-29-12)]. The methods in [\[7](#page-29-11)] and [\[9](#page-29-12)], however, differ significantly from those used in $[6]$. In fact, while the authors in $[7, 9]$ $[7, 9]$ $[7, 9]$ $[7, 9]$ work explicitly with the differential expressions and operators, the result in [\[6](#page-29-8)] follows from an abstract theorem on bounded perturbations of *J* -non-negative operators. The idea is simple: we write $A = A_0 + V$, where

$$
A_0: f \mapsto -\operatorname{sgn} \cdot f'', \qquad f \in \operatorname{dom} A_0 = H^2(\mathbb{R}), \tag{1.2}
$$

and *V* is the operator of multiplication with sgn $\cdot q$. If *J* further denotes multiplication with sgn, then $J = J^* = J^{-1}$ and A_0 is *J*-non-negative, i.e., JA_0 is self-adjoint and non-negative. Hence, the abstract theorem can be applied.

Here, we proceed in a similar way as in $[6]$ $[6]$: we first prove an abstract theorem on relatively bounded perturbations of *J* -non-negative operators and apply it to the above situation. In this case, we require $q \in L^p(\mathbb{R})$ with $p \in [2, \infty]$ since these potentials lead to relatively bounded perturbations of *A*0. As a consequence, we obtain bounds on the non-real eigenvalues of *A* (see Theorem [6.2\)](#page-23-0); these are of the same flavor as those in [\[7\]](#page-29-11), but surprisingly improve them significantly in the case $q < 0$ (see Fig. [5\)](#page-27-0).

Let us now touch upon the abstract situation in more detail. For this, let $J \neq I$ be a self-adjoint involution on a Hilbert space $(H, (\cdot, \cdot))$, i.e., $J = J^* = J^{-1}$. A linear operator *A* in H is called *J*-self-adjoint if the composition *JA* is a self-adjoint operator in *H*. In contrast to self-adjoint operators, the spectrum of a *J* -self-adjoint operator is not necessarily real. The only restriction, in general, is its symmetry with respect to the real axis and it is not hard to construct *J* -self-adjoint operators whose spectrum covers the entire complex plane (see, e.g., Example [5.4\)](#page-18-0) or is empty. It is therefore reasonable to consider special classes of *J* -self-adjoint operators—e.g., (locally) definitizable operators [\[26](#page-29-13), [34\]](#page-30-14), or, more specific, *J* -non-negative operators. A *J* -self-adjoint operator *A* with non-empty resolvent set is called *J* -non-negative if $JA \geq 0$. It is well known [\[34\]](#page-30-14) that a *J*-non-negative operator has real spectrum and possesses a spectral function on $\overline{\mathbb{R}}$ with the possible singularities 0 and ∞ .

Several types of perturbations of *J* -non-negative operators have already been investigated in numerous works. Here, we only mention [\[1,](#page-29-14) [4,](#page-29-15) [23,](#page-29-16) [27\]](#page-29-17) for compact and finite-rank perturbations, [\[2,](#page-29-18) [3,](#page-29-6) [6](#page-29-8), [22\]](#page-29-19) for bounded perturbations and [\[25](#page-29-20)] for form-bounded perturbations. The main result in [\[6](#page-29-8)] states that a bounded perturbation $A_0 + V$

of a *J*-non-negative operator A_0 which satisfies a certain regularity condition can be spectrally decomposed into a bounded operator and a *J* -non-negative operator. In particular, the non-real spectrum of the perturbed operator is contained in a compact set *K* and bounds on *K* were given explicitly. We call such operators *J* -non-negative over $\overline{\mathbb{C}}$ \ *K* (cf. Definition [2.2\)](#page-5-0).

Here, our focus lies on *relatively bounded* perturbations of the form $A_0 + V$, where *A*⁰ is *J* -non-negative and *V* is both *J* -symmetric and *A*0-bounded. Our main result, Theorem [5.2,](#page-18-1) generalizes that of $[6]$ $[6]$ and shows that whenever the A_0 -bound of *V* is sufficiently small, then the operator $A_0 + V$ is *J*-non-negative over $\mathbb{C} \setminus K$ with a compact set *K* which is specified in terms of the relative bounds of *V*.

A brief outline of the paper is as follows. First of all, we provide the necessary notions and definitions in Sect. [2.](#page-4-0) In Sect. [3](#page-6-0) we consider a self-adjoint operator *S* in a Hilbert space and an *S*-bounded operator *T* with *S*-bound less than one and provide enclosures for the set

$$
K_S(T) := \{ \lambda \in \varrho(S) : \|T(S - \lambda)^{-1}\| \ge 1 \}.
$$

As a by-product, this result yields an improvement of spectral enclosures from [\[10\]](#page-29-21) for the operator $S + T$ (cf. Corollary [3.5](#page-8-0) and Remark [3.6\)](#page-9-0). However, the main reason for considering the set $K_S(T)$ is Theorem [4.1](#page-11-0) in Sect. [4](#page-11-1) which states in particular that the non-real spectrum of *J* -self-adjoint diagonally dominant block operator matrices of the form

$$
S = \begin{pmatrix} S_+ & M \\ -M^* & S_- \end{pmatrix} \tag{1.3}
$$

is contained in $K_{S_+}(M) \cap K_{S_+}(M^*)$ if both the *S*_−-bound of *M* and the *S*₊-bound of *M*[∗] are less than one. It is now a key observation that our main object of investigation – the operator $A_0 + V$ – can be written in the form [1.3.](#page-3-0) A renormalization procedure then allows us to derive the main result, Theorem [5.2.](#page-18-1) In Sect. [6](#page-21-0) we apply Theorem [5.2](#page-18-1) to singular indefinite Sturm-Liouville operators.

Notation: Throughout, *H* and *K* stand for Hilbert spaces. In this paper, whenever we write $T : \mathcal{H} \to \mathcal{K}$, we mean that *T* is a linear operator mapping from dom $T \subset \mathcal{H}$ to K , where dom *T* does not necessarily coincide with H nor do we assume that it is dense in *H*. The space of all bounded operators $T : \mathcal{H} \to \mathcal{K}$ with dom $T = \mathcal{H}$ is denoted by $L(\mathcal{H}, \mathcal{K})$. As usual, we set $L(\mathcal{H}) := L(\mathcal{H}, \mathcal{H})$. The spectrum of an operator *T* in \mathcal{H} is denoted by $\sigma(T)$, its resolvent set by $\rho(T) = \mathbb{C} \setminus \sigma(T)$. The set of eigenvalues of *T* is called the *point spectrum* of *T* and is denoted by $\sigma_p(T)$. The *approximate point spectrum* $\sigma_{ap}(T)$ of *T* is the set of all $\lambda \in \mathbb{C}$ for which there exists a sequence $(f_n) \subset$ dom *T* such that $||f_n|| = 1$ for all $n \in \mathbb{N}$ and $(T - \lambda) f_n \to 0$ as $n \to \infty$. Throughout, for a set $\Delta \subset \mathbb{C}$ and $r > 0$ we use the notation $B_r(\Delta) := \{z \in \mathbb{C} : \text{dist}(z, \Delta) \leq r\}.$ For $z_0 \in \mathbb{C}$ we also write $B_r(z_0) := B_r(\{z_0\}) = \{z \in \mathbb{C} : |z - z_0| \le r\}$. These sets are intentionally defined to be closed.

2 Preliminaries

Let $J \in L(\mathcal{H})$ be a self-adjoint involution, i.e., $J = J^* = J^{-1}$. The operator *J* induces an (in general indefinite) inner product $[\cdot, \cdot]$ on \mathcal{H} :

$$
[f, g] = (Jf, g), \quad f, g \in \mathcal{H}.
$$

An operator *A* in *H* is called *J -self-adjoint* (*J -symmetric*) if *J A* is self-adjoint (symmetric, resp.). Equivalently, *A* is self-adjoint (resp. symmetric) with respect to the inner product $[\cdot, \cdot]$. The spectrum of a *J*-self-adjoint operator is symmetric with respect to R, that is, $\lambda \in \sigma(A)$ \Leftrightarrow $\overline{\lambda} \in \sigma(A)$. However, more cannot be said, in general. It is easy to construct examples of *J* -self-adjoint operators whose spectrum is the entire complex plane (see, e.g., Example [5.4](#page-18-0) below). Therefore, the literature usually focusses on special classes of *J* -self-adjoint operators or on local spectral properties as the spectral points of positive and negative type which we shall explain next.

Let *A* be a *J*-self-adjoint operator. The subset $\sigma_{+}(A)$ ($\sigma_{-}(A)$) of $\sigma(A)$ consists of the points $\lambda \in \sigma_{ap}(A)$ for which each sequence $(f_n) \subset \text{dom } A$ with $||f_n|| = 1$ for all $n \in \mathbb{N}$ and $(A - \lambda) f_n \to 0$ as $n \to \infty$ satisfies

$$
\liminf_{n\to\infty} [f_n, f_n] > 0 \quad \left(\limsup_{n\to\infty} [f_n, f_n] < 0, \text{ resp.}\right).
$$

A point $\lambda \in \sigma_{+}(A)$ ($\sigma_{-}(A)$) is called a *spectral point of positive (negative) type* of *A* and a set $\Delta \subset \mathbb{C}$ is said to be of positive (negative) type with respect to *A* if $\Delta \cap \sigma(A) \subset \sigma_{+}(A)$ ($\Delta \cap \sigma(A) \subset \sigma_{-}(A)$, resp.). The notion of the spectral points of positive and negative type was introduced in [\[36\]](#page-30-15) (see also [\[33](#page-30-16)]). It is immediate that $\lambda \in \sigma_+(A)$ implies that for each $f \in \text{ker}(A - \lambda)$, $f \neq 0$, we have $[f, f] > 0$. Hence, ker($A - \lambda$) is a *J*-positive subspace. In fact, much more holds (see [\[36](#page-30-15)]): We have $\sigma_{+}(A) \subset \mathbb{R}$ and, if $\lambda \in \sigma_{+}(A)$, then

- (1) there exists an open (complex) neighborhood *U* of λ such that each point in *U* is either contained in $\varrho(A)$ or in $\sigma_{+}(A)$. In particular, $U \setminus \mathbb{R} \subset \varrho(A)$;
- (2) there exists a local spectral function¹ *E* on *U* ∩ R such that for each Borel set Δ with $\Delta \subset U \cap \mathbb{R}$ the projection $E(\Delta)$ is *J*-self-adjoint and $(E(\Delta)H, [\cdot, \cdot])$ is a Hilbert space.

Roughly speaking, the part of the operator *A* with spectrum in $U \cap \mathbb{R}$ is a self-adjoint operator in a Hilbert space. Similar statements hold for the spectral points of negative type.

Definition 2.1 A *J*-self-adjoint operator *A* in H is said to be *J*-non-negative if $\varrho(A) \neq \emptyset$ and $\sigma(JA) \subset [0,\infty)$, that is, $(JAf, f) \geq 0$ for $f \in \text{dom } A$ (equivalently, $[Af, f] \ge 0$ for $f \in \text{dom } A$). The operator A is said to be *uniformly J*-positive if it is *J*-non-negative and $0 \in \rho(A)$, i.e., $\sigma(JA) \subset (0, \infty)$.

¹ For a definition of this notion we refer to, e.g., $[44,$ Definition 2.2].

It is well known that the spectrum of a *J* -non-negative operator is real and that (0, ∞) ∩ σ (*A*) ⊂ σ ₊(*A*) as well as (−∞, 0) ∩ σ (*A*) ⊂ σ _−(*A*), see, e.g., [\[34](#page-30-14)]. Consequently, *A* possesses a spectral function *E* on $\overline{\mathbb{R}}$ with the possible singularities 0 and ∞ . The projection $E(\Delta)$ is then defined for all Borel sets $\Delta \subset \mathbb{R}$ for which 0 ∈/ ∂ and ∞ ∈/ ∂. The points 0 and ∞ are called the *critical points* of *A*. If both $\|E([\varepsilon, 1])\|$ and $\|E([-1, -\varepsilon])\|$ are bounded as $\varepsilon \downarrow 0$, the point 0 is said to be a *regular* critical point of A. In this case, the spectral projection $E(\Delta)$ also exists if $0 \in \partial \Delta$. A similar statement holds for the point ∞ . In the sequel, we agree on calling a *J*-non-negative operator *regular* if its critical points 0 and ∞ both are regular.

As was shown in $[6]$ $[6]$, the perturbation of a regular J -non-negative operator A_0 with a bounded *J*-self-adjoint operator *V* leads to a *J*-self-adjoint operator $A = A_0 + V$ whose non-real spectrum is bounded and for which there exist $r_{+} > 0$ such that (r_+,∞) is of positive type and $(-\infty, -r_-)$ is of negative type with respect to *A*. Hence, the perturbed operator exhibits the same good spectral properties as a *J* -non-negative operator in the exterior of a compact set. We call such an operator *J* -non-negative in a neighborhood of ∞ . The following definition makes this more precise. Here, for a set $\Delta \subset \mathbb{C}$ we define $\Delta^* := {\overline{\lambda} : \lambda \in \Delta}$. By \mathbb{C}^+ we denote the open upper complex half-plane. We also set $\mathbb{R}^+ := (0, \infty)$ and $\mathbb{R}^- := (-\infty, 0)$.

Definition 2.2 Let $K = K^* \subset \mathbb{C}$ be a compact set, $0 \in K$, such that $\mathbb{C}^+ \setminus K$ is simply connected. A *J*-self-adjoint operator *A* in *H* is said to be *J*-non-negative over $\mathbb{C} \setminus K$ if the following conditions are satisfied:

- (i) $\sigma(A) \setminus \mathbb{R} \subset K$.
- (ii) $(\sigma(A) \cap \mathbb{R}^{\pm}) \setminus K \subset \sigma_{+}(A)$.
- (iii) There exist *M* > 0 and a compact set *K'* $\supset K$ such that for $\lambda \in \mathbb{C} \setminus (\mathbb{R} \cup K')$ we have

$$
\|(A - \lambda)^{-1}\| \le M \, \frac{(1 + |\lambda|)^2}{|\operatorname{Im} \lambda|^2}.
$$
 (2.1)

The relation [\(2.1\)](#page-5-1) means that the growth of the resolvent of *A* at ∞ is of order at most 2. The order is 1 if the fraction in [\(2.1\)](#page-5-1) can be replaced by $|\text{Im }\lambda|^{-1}$.

Remark 2.3 Note that the notion of *J*-non-negativity over $\overline{\mathbb{C}} \setminus K$ does not depend on *J* explicitly, but only on the inner product $[\cdot, \cdot]$. That is, if $(\cdot, \cdot)_0$ is an equivalent Hilbert space scalar product on *H* such that $[\cdot, \cdot] = (J_0, \cdot)$ for some (\cdot, \cdot) ₀-self-adjoint involution J_0 , then an operator A in H is J-non-negative over $\mathbb{C} \setminus K$ in $(\mathcal{H}, \langle \cdot, \cdot \rangle)$ if and only if it is J_0 -non-negative over $\overline{\mathbb{C}} \setminus K$ in $(\mathcal{H}, (\cdot, \cdot)_0)$.

Due to (ii) a *J*-self-adjoint operator *A* that is *J*-non-negative over $\overline{C} \setminus K$ possesses a (local) spectral function *E* on $\overline{\mathbb{R}} \setminus K$ with a possible singularity at ∞ . We say that *A* is *regular at* ∞ if ∞ is not a singularity of *E*. By [\[6](#page-29-8), Thm. 2.6 and Prop. 2.3] this is the case if and only if there exists a uniformly *J*-positive operator *W* in H such that *W* dom *A* ⊂ dom *A*.

In this paper we investigate relatively bounded perturbations of regular *J* -nonnegative operators. Recall that an operator $T : \mathcal{H} \to \mathcal{K}_1$ is called *relatively bounded* with respect to an operator $S : \mathcal{H} \to \mathcal{K}_2$ (or simply *S-bounded*) if dom $S \subset \text{dom } T$ and

$$
||Tf||^2 \le a||f||^2 + b||Sf||^2, \qquad f \in \text{dom } S,
$$
 (2.2)

where $a, b \ge 0$. The infimum of all possible *b* in [\(2.2\)](#page-5-2) is called the *S-bound* of *T*. It is often convenient to assume that the *S*-bound of *T* is less than one. Then, if *S* is closed, also $S + T$ is closed (see [\[28,](#page-29-22) Thm. IV.1.1]) and if *S* is self-adjoint and *T* symmetric, the sum $S + T$ is self-adjoint. The latter statement is known as the Kato-Rellich theorem (see, e.g., [\[28,](#page-29-22) Thm. V.4.3]).

3 Perturbations of Self-Adjoint Operators

In this section, *S* always denotes a self-adjoint operator in a Hilbert space *H* and $T : \mathcal{H} \to \mathcal{K}$ is an *S*-bounded operator, where \mathcal{K} is another Hilbert space. We note that in this situation we have

$$
[T(S - \lambda)^{-1}]^* = \overline{(S - \overline{\lambda})^{-1}T^*} \in L(K, \mathcal{H}),
$$
\n(3.1)

whenever $\lambda \in \varrho(S)$. The following set will play a crucial role in our spectral estimates in the subsequent sections:

$$
K_S(T) := \left\{ \lambda \in \varrho(S) : \|T(S - \lambda)^{-1}\| \ge 1 \right\}.
$$

Although its proof is elementary it seems that the following result is new.

Lemma 3.1 *Let H and K be Hilbert spaces, let S be a self-adjoint operator in H, and let* $T : \mathcal{H} \to \mathcal{K}$ *be such that*

$$
||Tf||^2 \le a||f||^2 + b||Sf||^2
$$
, $f \in \text{dom } S \subset \text{dom } T$,

where a, b \geq 0*, b* < 1*. Then*

$$
K_S(T) \subset B_S(T) := \bigcup_{t \in \sigma(S)} B_{\sqrt{a+bt^2}}(t). \tag{3.2}
$$

For $\lambda \notin B_S(T)$ *we have*

$$
||T(S - \lambda)^{-1}|| \le \sup_{t \in \sigma(S)} \frac{\sqrt{a + bt^2}}{|\lambda - t|} < 1.
$$
 (3.3)

Proof Let $f \in \text{dom } S^2$. Then

$$
||Tf||^2 \le (af, f) + (bS^2 f, f) = ||(a + bS^2)^{1/2} f||^2.
$$

If $\lambda \in \varrho(S)$, this implies

$$
||T(S - \lambda)^{-1}|| \le ||(a + bS^2)^{1/2}(S - \lambda)^{-1}|| = \sup_{t \in \sigma(S)} \phi_{\lambda}(t),
$$

where $\phi_{\lambda}(t) := \frac{\sqrt{a+bt^2}}{|\lambda-t|}$. Therefore and due to $\lim_{t\to\pm\infty} \phi_{\lambda}(t) = \sqrt{b} < 1$, for $||T(S - t)||$ λ)⁻¹ || < 1 it is sufficient that $\phi_{\lambda}(t)$ < 1 for each *t* ∈ σ(*S*). But the latter is just equivalent to λ not being contained in the right-hand side of (3.2). equivalent to λ not being contained in the right-hand side of [\(3.2\)](#page-6-1).

Corollary 3.2 *Under the assumptions of* Lemma [3.1](#page-6-2) *we have*

$$
K_S(T) \subset \left\{ \lambda \in \mathbb{C} : (\operatorname{Im} \lambda)^2 \le a + \frac{b}{1-b} \left(\operatorname{Re} \lambda \right)^2 \right\}.
$$
 (3.4)

Proof We prove that $B_{\sqrt{a+bt^2}}(t)$ is contained in the right-hand side of [\(3.4\)](#page-7-0) for arbitrary *t* ∈ R. For this, let $\lambda = \alpha + i\beta \in B_{\sqrt{a+bt^2}}(t)$, i.e., $(\alpha - t)^2 + \beta^2 < a + bt^2$, which implies

$$
\beta^2 < a + 2\alpha t - \alpha^2 - (1 - b)t^2 \le a + \frac{b}{1 - b}\alpha^2.
$$

Indeed, the last inequality follows from $(\frac{\alpha}{\sqrt{1-b}} - t\sqrt{1-b})^2 \ge 0$.

Remark 3.3 Let *S* and *T* be as in Lemma [3.1](#page-6-2) and $\lambda \in \mathbb{C} \setminus \mathbb{R}$. Let us consider the global behaviour of the function

$$
\varphi_{\lambda}(t) := \frac{a + bt^2}{|t - \lambda|^2}.
$$
\n(3.5)

First of all, we always have $\varphi_{\lambda}(\pm \infty) = b$. If $b = 0$, then φ_{λ} has no local minima but a global maximum at $t = \text{Re }\lambda$. Let $b > 0$, Re $\lambda \neq 0$, and set

$$
m_{\lambda} := \frac{b|\lambda|^2 - a}{2b \cdot \text{Re}\,\lambda}.
$$
 (3.6)

Then we have $\varphi_{\lambda}(m_{\lambda}) = \varphi_{\lambda}(\pm \infty) = b$, and φ_{λ} has a global maximum and a global minimum at

$$
t_{\text{max}} = m_{\lambda} + \text{sgn}(\text{Re }\lambda) \sqrt{\frac{a}{b} + m_{\lambda}^2}
$$
 and $t_{\text{min}} = m_{\lambda} - \text{sgn}(\text{Re }\lambda) \sqrt{\frac{a}{b} + m_{\lambda}^2}$,

respectively. Let *b* > 0, Re $\lambda = 0$. If $(\text{Im }\lambda)^2 > \frac{a}{b}$, then φ_{λ} has no local maxima but a global minimum at zero. If $(\text{Im }\lambda)^2 < \frac{a}{b}$, then φ_λ has no local minima but a global maximum at zero. In case $(\text{Im }\lambda)^2 = \frac{a}{b}$, we have $\varphi_{\lambda} \equiv b$.

Corollary 3.4 *Let H and K be Hilbert spaces, let S be a self-adjoint operator in H which is bounded from above by* $\gamma \in \mathbb{R}$ *, and let* $T : \mathcal{H} \to \mathcal{K}$ *be such that*

$$
||Tf||^2 \le a||f||^2 + b||Sf||^2
$$
, $f \in \text{dom } S \subset \text{dom } T$,

where $a \geq 0$ *,* $b \in (0, 1)$ *<i>. Then*

$$
K_S(T) \subset \bigcup_{t \le \gamma} B_{\sqrt{a+bt^2}}(t). \tag{3.7}
$$

 $\int f|\lambda| > \gamma + \sqrt{\gamma^2 + \frac{a}{b}}$ and $\text{Re }\lambda \geq 0$ *, then*

$$
||T(S - \lambda)^{-1}||^2 \leq b.
$$

Proof We only have to prove the last claim, which follows from [\(3.3\)](#page-6-3) if the function φ_λ from [\(3.5\)](#page-7-2) does not exceed the value *b* on ($-\infty$, γ]. Hence, let $|\lambda| > \gamma + \sqrt{\gamma^2 + \frac{a}{b}}$ and Re $\lambda \ge 0$. Then also $|\lambda|^2 > 2\gamma |\lambda| + \frac{a}{b} \ge 2\gamma$ (Re λ) + $\frac{a}{b}$. Assume that Re $\lambda = 0$. Then $(\text{Im }\lambda)^2 > \frac{a}{b}$, and Remark [3.3](#page-7-3) implies that $\varphi_\lambda(t) \leq b$ for all $t \in \mathbb{R}$. Let Re $\lambda > 0$, Im $\lambda \neq 0$. Then $|\lambda|^2 > 2\gamma$ (Re λ) + $\frac{a}{b}$ is equivalent to $m_{\lambda} > \gamma$ with m_{λ} as defined in [\(3.6\)](#page-7-4). Note that $\varphi_{\lambda}(m_{\lambda}) = b$ and that $t_{\max} > m_{\lambda} > \gamma$. Hence $\varphi_{\lambda}(t) \leq b$ for all $t \leq \gamma$. For Re $\lambda > 0$ and Im $\lambda = 0$ the claim follows by continuity of $\mu \mapsto T(S - \mu)^{-1}$ on $\varrho(S)$, which is an immediate consequence of $T(S - \lambda)^{-1} - T(S - \mu)^{-1} = (\lambda - \mu)[T(S - \lambda)^{-1}](S - \mu)^{-1}$. □ μ) $[T(S - \lambda)^{-1}] (S - \mu)^{-1}$.

Figure [1](#page-8-1) illustrates the situation in Corollary 3.4. A similar result holds for the case where *S* is bounded from below. It follows from Corollary [3.4](#page-7-5) by considering −*S*.

Let us briefly consider the situation where $K = H$ and thus $T : H \rightarrow H$. If $\lambda \in \rho(S)$ and $||T(S - \lambda)^{-1}|| < 1$, then

$$
S + T - \lambda = \left[I + T(S - \lambda)^{-1} \right] (S - \lambda)
$$

implies that also $\lambda \in \varrho(S + T)$. By contraposition, $\sigma(S + T) \subset \sigma(S) \cup K_S(T)$. This immediately leads to the following corollary.

Corollary 3.5 *Let H be a Hilbert space, let S be a self-adjoint operator in H, and let* $T : \mathcal{H} \to \mathcal{H}$ *be such that*

$$
||Tf||^2 \le a||f||^2 + b||Sf||^2
$$
, $f \in \text{dom } S \subset \text{dom } T$,

where a, b > 0 *, b* < 1 *. Then*

$$
\sigma(S+T) \subset \bigcup_{t \in \sigma(S)} B_{\sqrt{a+bt^2}}(t) \subset \left\{ \lambda \in \mathbb{C} : (\text{Im }\lambda)^2 \le a + \frac{b}{1-b} (\text{Re }\lambda)^2 \right\}. \tag{3.8}
$$

Remark 3.6 Corollary [3.5](#page-8-0) improves and refines the first two parts of Theorem 2.1 in [\[10](#page-29-21)]. The general spectral inclusion in [\[10](#page-29-21)] is

$$
\sigma(S+T) \subset \left\{\lambda \in \mathbb{C} : (\operatorname{Im} \lambda)^2 \leq \frac{a+b(\operatorname{Re} \lambda)^2}{1-b}\right\}.
$$

In the case where $a > 0$ the spectral enclosure [\(3.8\)](#page-9-1) is obviously strictly sharper. However, both boundary curves have the same asymptotes so that the improvement of the spectral inclusion only takes effect for $| \text{Re } \lambda |$ 'not too large' (see Fig. [2](#page-9-2) below). In the second part of [\[10](#page-29-21), Thm. 2.1] it was proved that if (σ_-, σ_+) is a spectral gap of *S* such that $\sigma'_{-} := \sigma_{-} + \sqrt{a + b\sigma_{-}^{2}} < \sigma_{+} - \sqrt{a + b\sigma_{+}^{2}} =: \sigma'_{+}$, then {λ : σ' $-$ < Re λ < σ'_+ } $\subset \rho(S+T)$. This now follows immediately from the first enclosure in [\(3.8\)](#page-9-1).

Example 3.7 We consider the massless Dirac operator $H_0 = -i(\sigma_1 \partial_x + \sigma_2 \partial_y)$ on \mathbb{R}^2 with Hermitian matrices $\sigma_1, \sigma_2 \in \mathbb{C}^{2 \times 2}$ as in [\[10](#page-29-21), Section 5.1]. Its domain is given by dom $H_0 = H^1(\mathbb{R}^2)^2 \subset L^2(\mathbb{R}^2)^2$. As in [\[10\]](#page-29-21) one shows that for a potential

 $V \in L^p(\mathbb{R}^2, \mathbb{C}^{2 \times 2}), p > 2$, one has

$$
||Vf||_2^2 \leq C_p^{\frac{2p}{p-2}} b^{-\frac{2}{p-2}} ||f||_2^2 + b||H_0f||_2^2, \quad f \in \text{dom } H_0,
$$

for every *b* ∈ (0, 1), where $C_p = \frac{\|V\|_p}{[2\pi (p-2)]^{1/p}}$. By Corollary [3.5,](#page-8-0) we have

$$
\sigma(H_0+V)\subset \bigcap_{b\in(0,1)}\left\{\lambda\in\mathbb{C}: (\text{Im }\lambda)^2\leq C_p^{\frac{2p}{p-2}}b^{-\frac{2}{p-2}}+\frac{b}{1-b}(\text{Re }\lambda)^2\right\}.
$$

The envelope of this set in the first quadrant is the curve $\gamma(b) = (x(b), y(b)), b \in$ (0, 1), where

$$
x(b) = \left(\frac{C_p}{b}\right)^{\frac{p}{p-2}} \cdot \sqrt{\frac{2}{p-2}} (1-b^2) \quad \text{and} \quad y(b) = \left(\frac{C_p}{b}\right)^{\frac{2}{p-2}} \cdot C_p \cdot \sqrt{\frac{p-2b^2}{p-2}}.
$$

The spectrum of $H_0 + V$ is thus bounded by this curve. Figure [3](#page-10-0) shows the curve γ and the spectral inclusion from [\[10\]](#page-29-21).

4 Spectral Properties of Diagonally Dominant J-Self-Adjoint Operator Matrices

In this section we consider a Hilbert space $\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_-$ (where \oplus denotes the orthogonal sum of subspaces) and operator matrices of the form

$$
S = \begin{pmatrix} S_+ & M \\ -M^* & S_- \end{pmatrix}, \qquad \text{dom } S = \text{dom } S_+ \oplus \text{dom } S_-, \tag{4.1}
$$

where S_+ and S_- are self-adjoint operators in the Hilbert spaces H_+ and H_- , respectively, and *M* : $H_-\to H_+$ is *S*₋-bounded such that also M^* is *S*₊-bounded. In particular, dom *S*[−] ⊂ dom *M* and dom *S*⁺ ⊂ dom *M*∗. Such operator matrices are called *diagonally dominant* (cf. [\[47,](#page-30-18) Def. 2.2.1]). It is clear that the operator *S* is *J* -symmetric, where *J* is the involution operator

$$
J = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix},
$$

which induces the indefinite inner product

$$
\begin{bmatrix}\n\begin{pmatrix} f_+ \\ f_- \end{pmatrix}, \begin{pmatrix} g_+ \\ g_- \end{pmatrix}\n\end{bmatrix} = \begin{pmatrix} \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix} \begin{pmatrix} f_+ \\ f_- \end{pmatrix}, \begin{pmatrix} g_+ \\ g_- \end{pmatrix} \\
= (f_+, g_+) - (f_-, g_-), \quad f_\pm, g_\pm \in \mathcal{H}_\pm.
$$

The following theorem is essentially a generalization of [\[19,](#page-29-23) Thm. 4.3], where the operator *M* was assumed to be bounded.

Theorem 4.1 *Assume that both the S*−*-bound of M and the S*+*-bound of M*[∗] *are less than one. Then S is J -self-adjoint,*

$$
\sigma(S) \setminus \mathbb{R} \subset K_{S-}(M) \cap K_{S+}(M^*), \tag{4.2}
$$

and

(i) $(\mathbb{R} \cap \varrho(S_-)) \setminus K_{S_-}(M)$ *is of positive type with respect to S.*

(ii) $(\mathbb{R} \cap \varrho(S_+)) \setminus K_{S_+}(M^*)$ *is of negative type with respect to S.*

Moreover, if $\lambda \in \mathbb{C} \setminus \mathbb{R}$ *, setting*

$$
L_{+}(\lambda) := M^*(S_{+} - \lambda)^{-1}
$$
 and $L_{-}(\lambda) := M(S_{-} - \lambda)^{-1}$,

we have

$$
\|(S-\lambda)^{-1}\| \le \frac{1}{|\operatorname{Im}\lambda|} \cdot \begin{cases} \frac{1+||L-(\lambda)||+||L-(\lambda)||^2}{1-||L-(\lambda)||^2} & \text{if } ||L-(\lambda)|| < 1\\ \frac{1+||L+(\lambda)||+||L+(\lambda)||^2}{1-||L+(\lambda)||^2} & \text{if } ||L+(\lambda)|| < 1. \end{cases} \tag{4.3}
$$

Proof By assumption, there exist constants $a_+, b_+ \geq 0, b_+ < 1$, such that

$$
||Mf_{-}||^{2} \le a_{-}||f_{-}||^{2} + b_{-}||S_{-}f_{-}||^{2}, \quad f_{-} \in \text{dom } S_{-} \subset \text{dom } M
$$

and

$$
||M^* f_+||^2 \le a_+ ||f_+||^2 + b_+ ||S_+ f_+||^2, \qquad f_+ \in \text{dom } S_+ \subset \text{dom } M^*.
$$

Define the operators

$$
S_0 := \begin{pmatrix} S_+ & 0 \\ 0 & S_- \end{pmatrix} \quad \text{and} \quad T := \begin{pmatrix} 0 & M \\ -M^* & 0 \end{pmatrix}
$$

with dom $S_0 = \text{dom } S_+ \oplus \text{dom } S_-$ and dom $T = \text{dom } M^* \oplus \text{dom } M$. Obviously, S_0 is *J*-self-adjoint and *T* is *J*-symmetric. Also, for $f = f_{+} + f_{-}$ with $f_{\pm} \in \text{dom } S_{\pm}$ we have

$$
||JTf||^{2} = ||Mf_{-}||^{2} + ||M^{*}f_{+}||^{2} \le a_{-}||f_{-}||^{2} + b_{-}||S_{-}f_{-}||^{2}
$$

+ $a_{+}||f_{+}||^{2} + b_{+}||S_{+}f_{+}||^{2}$
 $\le max\{a_{+}, a_{-}\}||f||^{2} + max\{b_{+}, b_{-}\}||JS_{0}f||^{2}.$

Hence, JT is JS_0 -bounded with JS_0 -bound less than one. By the Kato-Rellich theorem, $J(S_0 + T)$ is self-adjoint, which means that $S = S_0 + T$ is *J*-self-adjoint.

For the proof of [\(4.2\)](#page-11-2) let $\lambda \in \mathbb{C} \setminus \mathbb{R}$ such that $||M(S_-\lambda)^{-1}|| < 1$ (i.e., $||L_-(\lambda)|| <$ 1). We have to show that $\lambda \in \rho(S)$. For this, we make use of the first Schur complement of *S* which, for $\mu \in \varrho(S_+)$, is defined by

$$
S_1(\mu) = S_+ - \mu + M(S_- - \mu)^{-1}M^*
$$
, dom $S_1(\mu) = \text{dom } S_+$.

We shall exploit the fact that $\lambda \in \varrho(S)$ follows if $S_1(\lambda)$ is boundedly invertible (see, e.g., [\[47](#page-30-18), Thm. 2.3.3]). Since $||M(S_−\lambda)^{-1}|| < 1$, the operator $M(S_−\lambda)^{-1}M^*$ is *S*₊-bounded with *S*₊-bound *b*₊ < 1. Hence, *S*₁(λ) is closed (see [\[28,](#page-29-22) Thm. IV.1.1]). Now, for $\mu \in \rho(S_+)$ we have (cf. [\(3.1\)](#page-6-4))

$$
S_1(\lambda) - S_1(\mu) = \mu - \lambda + M[(S_- - \lambda)^{-1} - (S_- - \mu)^{-1}]M^*
$$

= $\mu - \lambda + (\lambda - \mu)M(S_- - \lambda)^{-1}(S_- - \mu)^{-1}M^*$
= $(\lambda - \mu)(L_-(\lambda)L_-(\overline{\mu})^*) - I|_{dom S_+}.$

From this it easily follows that $S_1(\mu)$ is closed for every $\mu \in \varrho(S_+)$. We also have $S_1(\overline{\mu}) \subset S_1(\mu)^*$. Therefore, for $f_+ \in \text{dom } S_+$, $||f_+|| = 1$,

Im
$$
(S_1(\lambda) f_+, f_+) = \frac{1}{2i} ((S_1(\lambda) - S_1(\overline{\lambda})) f_+, f_+)
$$

= $(Im \lambda) ((L_-(\lambda) L_-(\lambda)^* - I) f_+, f_+).$

But $||L_-(\lambda)|| < 1$. Hence, if Im $\lambda > 0$ (for Im $\lambda < 0$ a similar argument applies), the numerical range $W(S_1(\lambda))$ of $S_1(\lambda)$ is contained in the half-plane $H := \{z :$ Im $z \le -\delta$, where $\delta := (1 - ||L_-(\lambda)||^2)(\text{Im }\lambda) > 0$. Hence, for every $\mu \notin H$ and $f_+ \in \text{dom } S_+, || f_+ || = 1,$

$$
\|(S_1(\lambda) - \mu)f_+\| \ge |(S_1(\lambda)f_+, f_+) - \mu| \ge \text{dist}(\mu, H) = \text{Im}\,\mu + \delta. \tag{4.4}
$$

This implies that for each μ with Im $\mu > -\delta$ the operator $S_1(\lambda) - \mu$ is injective and has closed range. Let us now prove that there exists $r > 0$ such that $ri \in \rho(S_1(\lambda))$. Then it follows (cf. [\[28](#page-29-22), Ch. IV.5.6]) that $0 \in \varrho(S_1(\lambda))$ and thus $\lambda \in \varrho(S)$. Set $\mu := \lambda + ri$. Then

$$
S_1(\lambda) - ri = S_+ - \mu + L_-(\lambda)M^* = (I + L_-(\lambda)M^*(S_+ - \mu)^{-1})(S_+ - \mu).
$$

By Lemma [3.1](#page-6-2) the bounded operator $M^*(S_+ - \mu)^{-1}$ has norm less than one for sufficiently large *r*. This shows that $S_1(\lambda) - r_i$ is boundedly invertible for such *r*.

Above we have concluded from $|M(S_ - - \lambda)^{-1}| < 1$ that $\lambda \in \rho(S)$. If $|M^*(S_ + \lambda$)⁻¹ || < 1, one obtains $\lambda \in \varrho(S)$ by means of analogous arguments, using the second Schur complement $S_2(\lambda) = S_ - - \lambda + M^*(S_+ - \lambda)^{-1}M$.

Let us now prove the estimate [\(4.3\)](#page-11-3) for the resolvent of *S*. For this, let $\lambda \in \mathbb{C} \setminus \mathbb{R}$ such that $||L_-(\lambda)|| < 1$ as above. By [\[47](#page-30-18), Thm. 2.3.3] we have

$$
(S - \lambda)^{-1} = \begin{pmatrix} I & 0 \\ L_{-}(\overline{\lambda})^* & I \end{pmatrix} \begin{pmatrix} S_1(\lambda)^{-1} & 0 \\ 0 & (S_{-} - \lambda)^{-1} \end{pmatrix} \begin{pmatrix} I & -L_{-}(\lambda) \\ 0 & I \end{pmatrix}.
$$
 (4.5)

Denote the last factor by *L*. Then

$$
||L^*L|| = \left\| \begin{pmatrix} I & 0 \\ -L_{-}(\lambda)^* I \end{pmatrix} \begin{pmatrix} I & -L_{-}(\lambda) \\ 0 & I \end{pmatrix} \right\| = \left\| \begin{pmatrix} I & -L_{-}(\lambda) \\ -L_{-}(\lambda)^* I + L_{-}(\lambda)^* L_{-}(\lambda) \end{pmatrix} \right\|
$$

\n
$$
\leq \left\| \begin{pmatrix} I & 0 \\ 0 & I \end{pmatrix} \right\| + \left\| \begin{pmatrix} 0 & 0 \\ 0 & L_{-}(\lambda)^* L_{-}(\lambda) \end{pmatrix} \right\|
$$

\n
$$
+ \left\| \begin{pmatrix} 0 & L_{-}(\lambda) \\ L_{-}(\lambda)^* & 0 \end{pmatrix} \right\| = 1 + ||L_{-}(\lambda)||^2 + ||L_{-}(\lambda)||.
$$

Note that, since $(S_ - - \lambda)^{-1}$ is normal, we have

$$
||L_{-}(\overline{\lambda})|| = ||M(S_{-} - \overline{\lambda})^{-1}|| = ||(S_{-} - \lambda)^{-1}M^{*}|| = ||(S_{-} - \overline{\lambda})^{-1}M^{*}||
$$

= $||M(S_{-} - \lambda)^{-1}|| = ||L_{-}(\lambda)||$.

Therefore, for the first factor in [\(4.5\)](#page-13-0) we have the same estimate as for the last. For the middle factor in [\(4.5\)](#page-13-0) we obtain from [\(4.4\)](#page-13-1) with $\mu = 0$ that

$$
||S_1(\lambda)^{-1}|| \leq \delta^{-1} = |\operatorname{Im} \lambda|^{-1} (1 - ||L_{-}(\lambda)||^2)^{-1}
$$

and therefore [\(4.3\)](#page-11-3) in the case $\|L_-(\lambda)\| < 1$ follows. The estimate for the other case can be derived similarly by using the second Schur complement.

In the following last step of this proof we shall show that (i) holds true. The proof of (ii) follows analogous lines. Let $\lambda \in \mathbb{R} \cap \sigma(S) \cap \rho(S_-)$ such that $\tau := M(S_- \lambda$)⁻¹ || < 1. We have to show that $\lambda \in \sigma_+(S)$. For this, assume that $(f_n) \subset \text{dom } S$ is a sequence with $||f_n|| = 1$ and $(S - \lambda)f_n \to 0$ as $n \to \infty$. Let $f_n = f_n^+ + f_n^-$ with $f_n^{\pm} \in \text{dom } S_{\pm}$. Then

$$
-M^*f_n^+ + (S_- - \lambda)f_n^- \to 0, \qquad n \to \infty,
$$

in *H*_−. Applying $-(S_ - - \lambda)^{-1}$ gives (cf. [\(3.1\)](#page-6-4))

$$
[M(S_{-}-\lambda)^{-1}]^* f_n^+ - f_n^- \to 0, \quad n \to \infty.
$$

Set $g_n^- := [M(S_--\lambda)^{-1}]^* f_n^+, n \in \mathbb{N}$. Since both (g_n^-) and (f_n^-) are bounded sequences, we conclude that $\varepsilon_n := \|g_n^-\|^2 - \|f_n^-\|^2 \to 0$ as $n \to \infty$ and so

$$
||f_n^+||^2 = 1 - ||f_n^-||^2 = 1 - ||g_n^-||^2 + \varepsilon_n \ge 1 - \tau^2 ||f_n^+||^2 + \varepsilon_n,
$$

that is, $|| f_n^+ ||^2 \ge \frac{1}{2(1+\tau^2)}$ for *n* sufficiently large. Therefore, we obtain

$$
[f_n, f_n] = ||f_n^+||^2 - ||f_n^-||^2 = ||f_n^+||^2 - ||g_n^-||^2 + \varepsilon_n \ge (1 - \tau^2) ||f_n^+||^2
$$

$$
+ \varepsilon_n \ge \frac{1 - \tau^2}{2(1 + \tau^2)} + \varepsilon_n
$$

and thus, indeed, $\lambda \in \sigma_+(S)$.

Although (4.2) in Theorem [4.1](#page-11-0) is a fairly accurate estimate on the non-real spectrum, it requires complete knowledge about the norms of $M(S_ - - \lambda)^{-1}$ and $M^*(S_ + - \lambda)^{-1}$ for all $\lambda \in \mathbb{C} \setminus \mathbb{R}$. However, by making use of Lemma [3.1](#page-6-2) we immediately obtain the following theorem, where the spectral inclusion is expressed in terms of the spectra of *S*[−] and *S*⁺ instead of parameter-dependent norms.

Theorem 4.2 *Consider the operator matrix S in* [\(4.1\)](#page-11-4)*with self-adjoint diagonal entries S*⁺ *and S*[−] *and assume that*

$$
||Mf_{-}||^{2} \le a_{-}||f_{-}||^{2} + b_{-}||S_{-}f_{-}||^{2}, \quad f_{-} \in \text{dom } S_{-} \subset \text{dom } M,
$$

and

$$
||M^* f_+||^2 \le a_+ ||f_+||^2 + b_+ ||S_+ f_+||^2, \quad f_+ \in \text{dom } S_+ \subset \text{dom } M^*,
$$

where $a_+, b_+ \geq 0, b_+ < 1$ *. Then*

$$
\sigma(S) \setminus \mathbb{R} \subset \left(\bigcup_{t \in \sigma(S_-)} B_{\sqrt{a_- + b_- t^2}}(t)\right) \cap \left(\bigcup_{t \in \sigma(S_+)} B_{\sqrt{a_+ + b_+ t^2}}(t)\right). \tag{4.6}
$$

Moreover, setting $\Delta_{\pm}(t) := [t - \sqrt{a_{\pm} + b_{\pm}t^2}, t + \sqrt{a_{\pm} + b_{\pm}t^2}]$ *, we have* (i) $\mathbb{R} \setminus \bigcup_{t \in \sigma(S_-)} \Delta_-(t)$ *is of positive type with respect to S.* (ii) $\mathbb{R} \setminus \bigcup_{t \in \sigma(S_+)} \Delta_+(t)$ *is of negative type with respect to S.*

Remark 4.3 If *M* is bounded (i.e., $b_{\pm} = 0$ and $a_{\pm} = ||M||^2$), Theorem [4.2](#page-14-0) implies that $\sigma(S) \setminus \mathbb{R} \subset B_{\|M\|}(\sigma(S_-)) \cap B_{\|M\|}(\sigma(S_+))$, that $\mathbb{R} \setminus B_{\|M\|}(\sigma(S_-))$ is of positive type and $\mathbb{R} \setminus B_{\|M\|}(\sigma(S_+))$ is of negative type with respect to *S*. This result was already obtained in [\[6,](#page-29-8) Thm. 3.5] and more general spectral enclosures for $\sigma(S) \setminus \mathbb{R}$ have recently been found in [\[19\]](#page-29-23). Analogous methods might also lead to more general enclosures in the relatively bounded case. However, to avoid technical details we shall not touch this topic here.

For a short discussion of [\(4.6\)](#page-14-1), assume for simplicity that $a := a_+ = a_-$ and $b := b_+ = b_-.$ Clearly, if, e.g., $\sigma(S_-) = \sigma(S_+) = \mathbb{R}$, the intersection in [\(4.6\)](#page-14-1) does not improve the unions. On the other hand, if, e.g., $S_$ is bounded from above by $\gamma_$ and S_+ is bounded from below by γ_+ , we obtain from [\(4.6\)](#page-14-1) that

$$
\sigma(S) \setminus \mathbb{R} \subset \left(\bigcup_{t \leq \gamma_-} B_{\sqrt{a+bt^2}}(t)\right) \cap \left(\bigcup_{t \geq \gamma_+} B_{\sqrt{a+bt^2}}(t)\right).
$$

That is,

$$
\sigma(S) \setminus \mathbb{R} \subset \bigcup_{t \in [\gamma_+, \gamma_-]} B_{\sqrt{a+bt^2}}(t)
$$

if $\gamma_+ \leq \gamma_-$ and

$$
\sigma(S) \setminus \mathbb{R} \subset B_{\sqrt{a+b\gamma_+^2}}(\gamma_+) \cap B_{\sqrt{a+b\gamma_-^2}}(\gamma_-)
$$

if $\gamma_+ > \gamma_-$. The following corollary treats the case where, in addition, $\gamma_- = \gamma$ and $\gamma_+ = -\gamma$ for some $\gamma \geq 0$, which becomes relevant in the next section.

Corollary 4.4 *Consider the operator matrix S in* [\(4.1\)](#page-11-4) *with* $S_+ \geq -\gamma$ *and* $S_- \leq \gamma$ *with some* $\gamma \geq 0$ *and assume that*

$$
||Mf_{-}||^{2} \le a||f_{-}||^{2} + b||S_{-}f_{-}||^{2}, \quad f_{-} \in \text{dom } S_{-} \subset \text{dom } M,
$$

and

$$
||M^* f_+||^2 \le a||f_+||^2 + b||S_+ f_+||^2, \quad f_+ \in \text{dom } S_+ \subset \text{dom } M^*,
$$

where a, b \geq 0, *b* < 1*. Then the operator S in [\(4.1\)](#page-11-4) is J-non-negative over* $\overline{C} \setminus K$, *where*

$$
K = \bigcup_{t \in [-\gamma, \gamma]} B_{\sqrt{a+bt^2}}(t),\tag{4.7}
$$

and is regular at ∞ *.*

The following Fig. [4](#page-16-0) illustrates the set K in equation (4.7) in a special case.

Fig. 4 The set K in (4.7) with $a = \gamma = 10$ and $b = 0.4$. The blue segment represents the interval $[-\gamma, \gamma]$ (Color figure online)

Proof of Corollary [4.4](#page-15-1) It follows directly from Theorem [4.2](#page-14-0) and the preceding discussion that the non-real spectrum of *S* is contained in *K*. Theorem [4.2](#page-14-0) also implies that $\mathbb{R}^+ \setminus K = (\gamma + \sqrt{a + b\gamma^2}, \infty)$ is of positive type and $\mathbb{R}^- \setminus K = (-\infty, -\gamma - \sqrt{a + b\gamma^2})$ is of negative type with respect to *S*. $\sqrt{a + b\gamma^2}$ is of negative type with respect to *S*.

Let us prove that the growth of the resolvent of *S* at ∞ is of order at most 2 (cf. Definition [2.2](#page-5-0) (iii)). In fact, we prove that the order is 1. For this, let $\lambda \in \mathbb{C} \setminus \mathbb{R}$ such that $|\lambda| > \gamma + \sqrt{\gamma^2 + \frac{a}{b}}$. If Re $\lambda \ge 0$, then by Corollary [3.4](#page-7-5) we have $||M(S_ - \lambda)^{-1}||^2 \le b$. Thus, from (4.3) in Theorem [4.1](#page-11-0) we obtain

$$
\|(S-\lambda)^{-1}\| \le \frac{3}{|\operatorname{Im} \lambda| \cdot (1 - \|M(S_{-} - \lambda)^{-1}\|^2)} \le \frac{3}{1-b} \cdot \frac{1}{|\operatorname{Im} \lambda|}.
$$

A similar reasoning applies to the case where $\text{Re }\lambda \leq 0$.

The regularity of *S* at ∞ follows from [\[6](#page-29-8), Prop. 2.3] (see also [\[11\]](#page-29-24)) since *J* is iformly *I*-nositive and leaves dom *S* invariant uniformly *J* -positive and leaves dom *S* invariant.

5 A Perturbation Result for *J***-Non-Negative Operators**

Let *J* be a self-adjoint involution in the Hilbert space $(\mathcal{H}, \langle \cdot, \cdot \rangle)$ inducing the inner product $[\cdot, \cdot] = (J \cdot, \cdot)$ and let A_0 be a *J*-non-negative operator in H with spectral function *E*. Assume that both 0 and ∞ are not singular critical points of A_0 (i.e., *A*₀ is regular) and $0 \notin \sigma_p(A_0)$. Then both spectral projections $E_{\pm} := E(\mathbb{R}^{\pm})$ exist, $E_{+} + E_{-} = I$, and $J_0 := E_{+} - E_{-}$ is a bounded and boundedly invertible *J*-nonnegative operator which satisfies $J_0^2 = I$. In particular, J_0 is both self-adjoint and unitary with respect to the (positive definite) scalar product $(\cdot, \cdot)_0$, where

$$
(f,g)_0 := [J_0f,g], \quad f,g \in \mathcal{H}.
$$

By $\|\cdot\|_0$ we denote the norm corresponding to $(\cdot, \cdot)_0$. Since $\|\cdot\|_0$ is equivalent to the original norm, $(\mathcal{H}, (\cdot, \cdot)_0)$ is a Hilbert space. According to [\[6](#page-29-8), Lemma 3.8] we have

$$
J_0 = \frac{1}{\pi} \cdot \operatorname*{s-lim}_{n \to \infty} \int_{1/n}^{n} \left((A_0 + it)^{-1} + (A_0 - it)^{-1} \right) dt,\tag{5.1}
$$

where s-lim stands for the strong limit. We set

$$
\tau_0 := \|J_0\|.\tag{5.2}
$$

Note that from $J_0^2 = I$ it follows that $\tau_0 \geq 1$. In what follows, by $T^{\mathcal{B}}$ we denote the adjoint of $T \in L(\mathcal{H})$ with respect to the scalar product $(\cdot, \cdot)_0$.

Lemma 5.1 *Let* $T \in L(H)$ *. Then*

$$
||T||_0 \leq \tau_0 ||T||.
$$

Moreover for $f \in H$ *, we have*

$$
\tau_0^{-1} \|f\|^2 \, \leq \|f\|_0^2 \, \leq \, \tau_0 \|f\|^2
$$

as well as

$$
\|E_{\pm}f\|_0^2 \le \frac{1+\tau_0}{2} \|f\|^2 \quad \text{and} \quad \|E_{\pm}f\|^2 \le \frac{1+\tau_0}{2} \|f\|_0^2.
$$

Proof Since the spectrum of a bounded operator is contained in the closure of its numerical range, we have

$$
\|T\|_{0}^{2} = \|T^{\circledast}T\|_{0} = \sup\{|\lambda| : \lambda \in \sigma(T^{\circledast}T)\} \le \sup\{|(T^{\circledast}Tf, f)| : f \in \mathcal{H}, \|f\| = 1\}
$$

$$
\le \|T^{\circledast}\| \|T\|.
$$

Hence, from $T^{\circledast} = J_0 J T^* J J_0$ we obtain $||T^{\circledast}|| \le ||J_0||^2 ||T|| = \tau_0^2 ||T||$.

Let now $f \in H$. Then we immediately obtain $||f||_0^2 = [J_0 f, f] \le ||J_0|| ||f||^2 =$ $\tau_0 || f ||^2$ and, by the first claim,

$$
||f||^2 = [f, Jf] = (J_0f, Jf)_0 \le ||J||_0 ||f||_0^2 \le \tau_0 ||f||_0^2.
$$

Moreover,

$$
||E_{\pm}f||_0^2 = (E_{\pm}f, f)_0 = \frac{1}{2}(f \pm J_0f, f)_0 = \frac{1}{2}[J_0f \pm f, f] = \frac{1}{2}(J_0f \pm f, Jf).
$$

Hence, $||E_{\pm} f||_0^2 \le \frac{1}{2} ||I \pm J_0 || ||f||^2 \le \frac{1+\tau_0}{2} ||f||^2$.

Let P_{\pm} denote the orthogonal projection onto $H_{\pm} := E_{\pm} \mathcal{H}$ with respect to the scalar product (·, ·). By [\[6,](#page-29-8) Lemma 3.9] we have $P_{\pm}J|_{\mathcal{H}_{\pm}} = (E_{\pm}J|_{\mathcal{H}_{+}})^{-1}$. Note that $P_+ J|_{\mathcal{H}_+}$ is a non-negative self-adjoint operator in $(\mathcal{H}_+, \overline{(\cdot, \cdot)})$. Hence, we have

$$
\min \sigma(P_+ J|_{\mathcal{H}_+}) = \left[\max \sigma(E_+ J|_{\mathcal{H}_+})\right]^{-1} = \|E_+ J|_{\mathcal{H}_+}\|^{-1} \ge \|E_+\|^{-1}
$$

$$
= \|\frac{1}{2}(I + J_0)\|^{-1} \ge \frac{2}{1 + \tau_0}.
$$

Therefore, for $f_+ \in \mathcal{H}_+$,

$$
|| f_{+}||_{0}^{2} = [f_{+}, f_{+}] = (Jf_{+}, f_{+}) = (P_{+}Jf_{+}, f_{+}) \ge \frac{2}{1+\tau_{0}}|| f_{+}||^{2},
$$

and consequently, $||E_+ f||^2 \le \frac{1+\tau_0}{2} ||E_+ f||_0^2 \le \frac{1+\tau_0}{2} ||f||_0^2$. A similar reasoning applies to $||E-f||^2$. $□$

We can now state and prove the main theorem in this section.

Theorem 5.2 *Let* A_0 *be a regular J-non-negative operator in* H *with* $0 \notin \sigma_p(A_0)$ *and let* $\tau > \tau_0$ *, where* τ_0 *is as in* [\(5.2\)](#page-17-0)*. Furthermore, let* V *be a J*-symmetric operator in *H with* dom *A*⁰ ⊂ dom *V such that*

$$
(1+\tau)\tau ||Vf||^2 \le 2a||f||^2 + b||A_0f||^2, \quad f \in \text{dom}\, A_0,
$$

where a, b \geq 0*, b* < 1*. Then the operator A := A₀ + V is J-self-adjoint.*

Let $v := \inf\{(JVf, f) : f \in \text{dom } V, ||f|| = 1\} \in [-\infty, \infty)$ *. If* $v \ge 0$ *, then A is a J -non-negative operator. Otherwise, the operator A is J -non-negative over*

$$
\overline{\mathbb{C}} \setminus \bigcup_{t \in [-\gamma, \gamma]} B_{\sqrt{a+bt^2}}(t), \tag{5.3}
$$

 $where^2 \gamma := min\{\sqrt{\frac{1+\tau}{2\tau}}a, -\frac{1+\tau}{2}v\}$ $where^2 \gamma := min\{\sqrt{\frac{1+\tau}{2\tau}}a, -\frac{1+\tau}{2}v\}$ $where^2 \gamma := min\{\sqrt{\frac{1+\tau}{2\tau}}a, -\frac{1+\tau}{2}v\}$ *. In both cases the operator A is regular at* ∞ *. If* $b < \frac{\tau - 1}{2\tau}$, then A is J-non-negative over

$$
\overline{\mathbb{C}} \setminus \bigcup_{t \in [-\gamma, \gamma]} B_{\sqrt{\frac{1+\tau}{2\tau(1-b)}} (a+bt^2)}(t). \tag{5.4}
$$

Remark 5.3 If $b < \frac{\tau-1}{2\tau}$, then the bound in [\(5.4\)](#page-18-3) is better than that in [\(5.3\)](#page-18-4). If *V* is bounded (i.e., $b = 0$ and $a = \frac{(1+\tau)\tau}{2} ||V||^2$), we obtain from the second part of Theorem [5.2](#page-18-1) that $A_0 + V$ is *J*-non-negative over

$$
\overline{\mathbb{C}}\setminus B_r([-d,d]),
$$

where $r = \frac{1+\tau}{2} ||V||$ and $d = -\frac{1+\tau}{2}$ min $\sigma(JV)$, which is the same result as [\[6,](#page-29-8) Thm. 3.1].

The following example (cf. [\[6,](#page-29-8) Example 3.2]) shows that the assumption on the regularity of A_0 in Theorem [5.2](#page-18-1) cannot be dropped.

Example 5.4 Let $(K, \langle \cdot, \cdot \rangle)$ be a Hilbert space and let *H* be an unbounded selfadjoint operator in *K* such that $\sigma(H) \subset (0, \infty)$.

² If $v = -\infty$, we set $\gamma = \sqrt{\frac{1+\tau}{2\tau}}a$.

Consider the following operators in $H := K \oplus K$:

$$
J = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}, \quad A_0 = \begin{pmatrix} 0 & I \\ H & 0 \end{pmatrix}, \quad V = \begin{pmatrix} 0 & -I \\ 0 & 0 \end{pmatrix}, \quad \text{and} \quad A_0 + V = \begin{pmatrix} 0 & 0 \\ H & 0 \end{pmatrix}.
$$

It is easy to see that A_0 is a *J*-non-negative operator and *V* is a bounded *J*-selfadjoint operator. Moreover, as dom $(A_0 + V) =$ dom $A_0 =$ dom $H \oplus K$ we conclude ran $(A_0 +$ $V - \lambda$) \neq *H* for every $\lambda \in \mathbb{C}$, that is, $\sigma(A_0 + V) = \mathbb{C}$.

Proof of Theorem [5.2](#page-18-1) Without loss of generality we assume that dom $V = \text{dom } A_0$. As $\tau \geq \tau_0 \geq 1$, it follows from the Kato-Rellich theorem that $A = A_0 + V$ is *J*-self-adjoint. Equivalently, *A* is J_0 -self-adjoint in $(\mathcal{H}, (\cdot, \cdot)_0)$.

Let \oplus_0 denote the $(\cdot, \cdot)_0$ -orthogonal sum of subspaces. Then we have $\mathcal{H} = \mathcal{H}_+ \oplus_0$ \mathcal{H}_- and dom $A_0 = D_+ \oplus_0 D_-,$ where $\mathcal{H}_\pm := (E_\pm \mathcal{H}, \langle \cdot, \cdot \rangle_0)$ and $D_\pm = \text{dom } A_0 \cap \mathcal{H}_\pm.$ Hence, *V* is defined on both D_+ and D_- and we can write A_0 and *V* as operator matrices

$$
A_0 = \begin{pmatrix} A_+ & 0 \\ 0 & A_- \end{pmatrix} \quad \text{and} \quad V = \begin{pmatrix} V_+ & M \\ N & V_- \end{pmatrix},
$$

where

$$
A_{\pm} = E_{\pm}(A|_{D_{\pm}}), \quad V_{\pm} = E_{\pm}(V|_{D_{\pm}}), \quad M = E_{+}(V|_{D_{-}}), \quad \text{and} \quad N = E_{-}(V|_{D_{+}}).
$$

Note that A_0 is J_0 -non-negative in $(\mathcal{H}, (\cdot, \cdot)_0)$ which implies that $\pm A_{\pm} \geq 0$ in \mathcal{H}_{\pm} . Moreover, since *V* is *J*-symmetric (and hence J_0 -symmetric in $(\mathcal{H}, (\cdot, \cdot)_0)$), we have that $N \subset -M^{\circledast}$. Hence,

$$
A = \begin{pmatrix} A_+ + V_+ & M \\ -M^{\otimes} & A_- + V_- \end{pmatrix}, \quad \text{dom } A = D_+ \oplus_0 D_-.
$$

In particular, $A_{\pm} + V_{\pm}$ is a self-adjoint operator in \mathcal{H}_{\pm} . For $f_{\pm} \in D_{\pm}$, we obtain from Lemma [5.1](#page-17-1) that

$$
\begin{aligned} \|V_{\pm}f_{\pm}\|_{0}^{2} &= \|E_{\pm}Vf_{\pm}\|_{0}^{2} \le \frac{1+\tau_{0}}{2} \|Vf_{\pm}\|^{2} \le \frac{a}{\tau} \|f_{\pm}\|^{2} + \frac{b}{2\tau} \|A_{\pm}f_{\pm}\|^{2} \\ &\le \frac{1+\tau}{2\tau} \big(a \|f_{\pm}\|_{0}^{2} + \frac{b}{2} \|A_{\pm}f_{\pm}\|_{0}^{2} \big). \end{aligned}
$$

Hence, V_{\pm} is A_{\pm} -bounded (in \mathcal{H}_{\pm}) with bounds $\frac{1+\tau}{2\tau}a$ and $\frac{1+\tau}{4\tau}b$. As $\frac{1+\tau}{4\tau}b < 1$, $\sqrt{\frac{1+\tau}{2\tau}}a$ and $A_+ + V_+$ is bounded from below by $-\gamma_1$ (both with respect to $(\cdot, \cdot)_0$). Corollary [3.5](#page-8-0) implies that the operator $A_+ + V_-$ is bounded from above by $\gamma_1 :=$ Now, for *f*[−] ∈ *D*[−] we have

$$
\|Mf_{-}\|_{0}^{2} + \|V_{-}f_{-}\|_{0}^{2} = \|Vf_{-}\|_{0}^{2} \le \tau_{0} \|Vf_{-}\|^{2} \le \frac{2}{1+\tau} \left(a\|f_{-}\|^{2} + \frac{b}{2}\|A_{-}f_{-}\|^{2}\right)
$$

\n
$$
\le a\|f_{-}\|_{0}^{2} + \frac{b}{2}\|A_{-}f_{-}\|_{0}^{2}
$$

\n
$$
\le a\|f_{-}\|_{0}^{2} + b \cdot \left(\|(A_{-} + V_{-})f_{-}\|_{0}^{2} + \|V_{-}f_{-}\|_{0}^{2}\right).
$$

Since $b < 1$, it follows that

$$
||Mf-||_0^2 \le a||f-||_0^2 + b||(A_- + V_-)f-||_0^2.
$$

Similarly, for $f_+ \in D_+$ one obtains

$$
||M^{\circledast} f_{+}||_{0}^{2} \leq a||f_{+}||_{0}^{2} + b||(A_{+} + V_{+})f_{+}||_{0}^{2}.
$$

As seen above, $A_+ + V_+$ is bounded from below by $-\gamma_1$ and $A_- + V_-$ is bounded from above by γ_1 . Corollary [4.4](#page-15-1) and Remark [2.3](#page-5-3) thus imply that the operator A is *J*-non-negative over $\overline{\mathbb{C}} \setminus K_1$, where

$$
K_1 = \bigcup_{t \in [-\gamma_1, \gamma_1]} B_{\sqrt{a+bt^2}}(t).
$$

In particular, $\rho(A) \neq \emptyset$. Hence, if the lower bound v of *JV* is non-negative, the operator *A* is *J*-non-negative. Assume that $v < 0$. Then for $f_+ \in D_+$ we have

$$
((A_{+} + V_{+})f_{+}, f_{+})_{0} \geq (V_{+}f_{+}, f_{+})_{0} = (Vf_{+}, E_{+}f_{+})_{0} = (Vf_{+}, J_{0}f_{+})_{0}
$$

=
$$
[Vf_{+}, f_{+}] = (JVf_{+}, f_{+}) \geq v \|f_{+}\|^{2} \geq \frac{1+\tau}{2} v \|f_{+}\|_{0}^{2},
$$

and similarly $((A_+ + V_-)f_-, f_-)_0 \le -\frac{1+\tau}{2}v \|f_-\|_0^2$ for $f_- \in D_-$. Thus, A is also *J*-non-negative over $\overline{\mathbb{C}} \setminus K_2$, where

$$
K_2 = \bigcup_{t \in [-\gamma_2, \gamma_2]} B_{\sqrt{a+bt^2}}(t)
$$

and $\gamma_2 = -\frac{1+\tau}{2}v$. This finishes the proof of the first part of the theorem. Since *A*⁰ is regular at ∞ , there exists a uniformly *J*-positive operator *W* in *H* such that *W* dom $A_0 \subset$ dom A_0 (see [\[11\]](#page-29-24)). From dom $A =$ dom A_0 it follows that A is regular at ∞ .

For the second part of the theorem, let *f*[−] ∈ *D*−. Then we have

$$
||V_{-}f_{-}||^{2} = ||E_{-}Vf_{-}||^{2} \le \frac{1+\tau}{2}||Vf_{-}||_{0}^{2} \le \frac{(1+\tau)\tau}{2}||Vf_{-}||^{2}.
$$

Therefore,

$$
\frac{1+\tau}{2} ||Vf_{-}||^{2} \leq \frac{a}{\tau} ||f_{-}||^{2} + \frac{b}{2\tau} ||A_{-}f_{-}||^{2}
$$
\n
$$
\leq \frac{a}{\tau} ||f_{-}||^{2} + \frac{b}{\tau} ||(A_{-} + V_{-})f_{-}||^{2} + \frac{b}{\tau} ||V_{-}f_{-}||^{2}
$$
\n
$$
\leq \left(\frac{(1+\tau)a}{2\tau} ||f_{-}||_{0}^{2} + \frac{(1+\tau)b}{2\tau} ||(A_{-} + V_{-})f_{-}||_{0}^{2}\right) + \frac{1+\tau}{2}b||Vf_{-}||^{2},
$$

which implies

$$
||Mf-||_0^2 = ||E+Vf-||_0^2 \le \frac{1+\tau}{2} ||Vf-||^2
$$

$$
\leq \frac{(1+\tau)a}{2\tau(1-b)} \|f_-\|_0^2 + \frac{(1+\tau)b}{2\tau(1-b)} \|(A_- + V_-)f_-\|_0^2.
$$

Similarly, for $f_+ \in D_+$ we have

$$
||M^{\mathfrak{B}} f_{+}||_{0}^{2} \leq \frac{(1+\tau)a}{2\tau(1-b)}||f_{+}||_{0}^{2} + \frac{(1+\tau)b}{2\tau(1-b)}||(A_{+} + V_{+})f_{+}||_{0}^{2}.
$$

We apply Corollary [4.4](#page-15-1) and conclude that the claim holds if only $\frac{(1+\tau)b}{2\tau(1-b)} < 1$, which is equivalent to *b* < $\frac{2\tau}{1+3\tau}$. But since the first part of the theorem yields better bounds for $\frac{\tau-1}{2}$ < *b* < $\frac{2\tau}{1+2}$ (e.g., *b* = $\frac{1}{2}$), we require that *b* < $\frac{\tau-1}{2}$ for $\frac{\tau-1}{2\tau} \le b < \frac{2\tau}{1+3\tau}$ (e.g., $b = \frac{1}{2}$), we require that $b < \frac{\tau-1}{2\tau}$ $\frac{1}{2\tau}$

Since the *A*₀-bound of an *A*₀-compact operator is always zero (see [\[15,](#page-29-25) Corollary III.7.7]), the next corollary immediately follows from Theorem [5.2.](#page-18-1)

Corollary 5.5 *Let* A_0 *be a regular J-non-negative operator in* H *with* $0 \notin \sigma_p(A_0)$ *and let V be a J-symmetric operator in* H *which is A*⁰-*compact. Then the operator* $A_0 + V$ *is J-self-adjoint, J-non-negative over* $\overline{\mathbb{C}} \setminus K$ *for some K, and regular at* ∞ *.*

6 An Application to Singular Indefinite Sturm-Liouville Operators

In this section we apply the abstract Theorem [5.2](#page-18-1) to indefinite Sturm-Liouville operators arising from differential expressions of the form

$$
\mathcal{L}(f)(x) = \text{sgn}(x)\big(-f''(x) + q(x)f(x)\big), \quad x \in \mathbb{R},
$$

with a real-valued potential $q \in L^p(\mathbb{R})$, $2 \leq p \leq \infty$. The maximal domain corresponding to $\mathcal L$ is given by

$$
D_{\max} = \{ f \in L^{2}(\mathbb{R}) : f, f' \in AC_{\text{loc}}(\mathbb{R}), -f'' + qf \in L^{2}(\mathbb{R}) \},\
$$

where $AC_{loc}(\mathbb{R})$ stands for the space of functions $f : \mathbb{R} \to \mathbb{C}$ which are locally absolutely continuous. The maximal operator *A* associated with *L* is defined as

$$
Af := \mathcal{L}(f), \qquad \text{dom } A := D_{\text{max}}.\tag{6.1}
$$

We also define the corresponding *definite* Sturm-Liouville operator *T* by

$$
Tf := JAf = -f'' + qf, \qquad \text{dom } T := D_{\text{max}},
$$

where $J: L^2(\mathbb{R}) \to L^2(\mathbb{R})$ is the operator of multiplication with the function sgn(*x*). Note that $J = J^{-1} = J^*$ and hence *J* is a fundamental symmetry.

If $q \in L^p(\mathbb{R})$, $2 \le p \le \infty$, it is well known that the operator *T* is self-adjoint and bounded from below. In fact, this holds for much more general differential expressions (cf. [\[7](#page-29-11), Thm. 1.1]). Consequently, the operator $A = JT$ is *J*-self-adjoint and it was shown in [\[5,](#page-29-7) Thm. 4.2] that *A* is non-negative over $\mathbb{C}\setminus K$ for some compact $K = K^*$. In particular, the non-real spectrum of *A* is contained in *K*. It is also known (cf. [\[5](#page-29-7)]) that the non-real spectrum of *A* consists of isolated eigenvalues which can only accummulate to $[\mu, -\mu]$, where μ is the infimum of the essential spectrum of *T*. However, the set K could not explicitly be specified in [\[5\]](#page-29-7).

Recently, in [\[7\]](#page-29-11) spectral enclosures for the non-real spectrum have been found: $\sigma(A) \setminus \mathbb{R} \subset K'$. However, this does not mean that the operator *A* is non-negative over $\mathbb{C} \setminus K'$ for K' might be much smaller than K. In Theorem [6.2](#page-23-0) below, we provide explicit bounds on the compact set K, thereby improving the bounds from [\[7\]](#page-29-11) if $q < 0$.

The following lemma sets the ground for the applicability of Theorem [5.2](#page-18-1) to the problem. In its proof we make use of the Fourier transform. Here, we use the definition

$$
\mathcal{F}f(\omega) = \hat{f}(\omega) = \int_{\mathbb{R}} f(x)e^{-2\pi ix\omega} dx, \quad \omega \in \mathbb{R},
$$

for Schwartz functions f . The Fourier transform $\mathcal F$ is extended to a unitary operator in $L^2(\mathbb{R})$ in the usual way. By $H^m(\mathbb{R})$ we denote the Sobolev space of regularity order $m > 0$ on \mathbb{R} , i.e.,

$$
H^{m}(\mathbb{R}) = \{ f \in L^{2}(\mathbb{R}) : (1 + \omega^{2})^{m/2} \hat{f} \in L^{2}(\mathbb{R}) \}.
$$

Recall that for $f \in H^1(\mathbb{R})$ we have $\hat{f}'(\omega) = 2\pi i \omega \hat{f}(\omega)$ and thus, for $f \in H^2(\mathbb{R})$, $\widehat{f''}(\omega) = -4\pi^2\omega^2\,\hat{f}(\omega).$

Lemma 6.1 *Let* $p \in [2, \infty]$ *. Then for any* $g \in L^p(\mathbb{R})$ *and any* $f \in H^2(\mathbb{R})$ *we have* $f g \in L^2(\mathbb{R})$ *and for any r* > 0 *the following inequality holds:*

$$
||fg||_2 \le (2r)^{1/p} \left(||f||_2 + \frac{1}{2\sqrt{3}\pi^2 pr^2} ||f''||_2 \right) ||g||_p.
$$

Proof We first observe that for $f \in H^2(\mathbb{R})$ we have $\omega^2 \hat{f} \in L^2(\mathbb{R})$ and hence, for $r > 0$,

$$
\|\hat{f}\|_{1} = \int_{-r}^{r} |\hat{f}(\omega)| d\omega + \int_{|\omega|>r} (4\pi^{2}\omega^{2})^{-1} \cdot |4\pi^{2}\omega^{2}\hat{f}(\omega)| d\omega
$$

\n
$$
\leq \sqrt{2r} \|\hat{f}\|_{2} + \left(2 \int_{r}^{\infty} \frac{d\omega}{16\pi^{4}\omega^{4}}\right)^{1/2} \left(\int_{|\omega|>r} |4\pi^{2}\omega^{2}\hat{f}(\omega)|^{2} d\omega\right)^{1/2} (6.2)
$$

\n
$$
\leq \sqrt{2r} \||f\|_{2} + \frac{\sqrt{2}}{4\sqrt{3}\pi^{2}} r^{-3/2} \|f''\|_{2} = \sqrt{2r} \ C_{r}(f),
$$

where $C_r(f) := ||f||_2 + \frac{1}{4\sqrt{3}\pi^2 r^2} ||f''||_2$. Thus, for $g \in L^2(\mathbb{R})$ we obtain

$$
\|gf\|_2 \le \|f\|_{\infty} \|g\|_2 \le \|f\|_1 \|g\|_2,
$$

which yields the desired inequality for $p = 2$. For $p = \infty$ the inequality is evident.

Now, let $p \in (2, \infty)$ and fix $f \in H^2(\mathbb{R})$. Set $T_f g := fg$. Then T_f is bounded both as an operator from $L^{\infty}(\mathbb{R})$ to $L^2(\mathbb{R})$ (with norm $|| f ||_2$) and from $L^2(\mathbb{R})$ to $L^2(\mathbb{R})$ (with norm $\leq \sqrt{2r} C_r(f)$). By the Riesz-Thorin interpolation theorem (see, e.g., [\[16,](#page-29-26) Thm. 6.27]), T_f is also bounded as an operator from $L^{2/\theta}(\mathbb{R})$ to $L^2(\mathbb{R})$, $\theta \in (0, 1)$, with norm

$$
\|T_f\|_{L^{2/\theta}(\mathbb{R})\to L^2(\mathbb{R})} \le \|f\|_2^{1-\theta} \left((2r)^{1/2}C_r(f) \right)^{\theta} = (2r)^{\theta/2} \|f\|_2 \left(1 + \frac{\|f\|_2^{-1} \|f''\|_2}{4\sqrt{3}\pi^2 r^2} \right)^{\theta}.
$$

Now, the claim follows from setting $p = 2/\theta$ and the simple inequality $(1 + c)^{\theta} \le$ $1 + c\theta$, which holds for $\theta \in (0, 1)$ and $c > 0$.

Theorem 6.2 *Let* $q \in L^p(\mathbb{R})$ *,* $p \in [2, \infty)$ *, and set*

$$
s_p := 4 - 3\sqrt{2} - 5p + 4\sqrt{2}p + \sqrt{44 - 31\sqrt{2} - 88p + 62\sqrt{2}p + 57p^2 - 40\sqrt{2}p^2}
$$

$$
f(s) := \sqrt{2 \frac{(17 + 12\sqrt{2})s + 4 + 3\sqrt{2}}{(3 + 2\sqrt{2})s - 1 - \sqrt{2}}},
$$

as well as

$$
C_p := (1 + \sqrt{2})\sqrt{3 - 2\sqrt{2}}\sqrt{\frac{2p}{2p-1}} \left(\frac{16\sqrt{2}(3 + 2\sqrt{2})^2}{3\pi^4 p} s_p\right)^{\frac{1}{4p-2}}.
$$

Then the singular indefinite Sturm-Liouville operator A from [\(6.1\)](#page-21-1) *with potential q is J*-non-negative over $\overline{C} \setminus K$, where

$$
K := \left\{ \lambda \in \mathbb{C} : |\operatorname{Im} \lambda| \le C_p f(s_p) \|q\|_p^{\frac{2p}{2p-1}}, \ | \operatorname{Re} \lambda| \le C_p \left(\sqrt{6 + 4\sqrt{2}} + f(s_p) \right) \|q\|_p^{\frac{2p}{2p-1}} \right\}.
$$
\n
$$
(6.3)
$$

Proof The differential operator A_0 , defined by

$$
A_0 f := -\operatorname{sgn} \cdot f'', \qquad \operatorname{dom} A_0 := H^2(\mathbb{R}),
$$

is *J*-self-adjoint. In fact, A_0 is *J*-non-negative and neither 0 nor ∞ is a singular critical point of A_0 , see [\[13\]](#page-29-10). Also, obviously, $0 \notin \sigma_p(A_0)$.

Step 1. (*Calculation of* τ_0) We have $\tau_0 = ||J_0||$, where (see [\(5.1\)](#page-16-1))

$$
J_0 = \frac{1}{\pi} \cdot \operatorname*{s-lim}_{n \to \infty} \int_{1/n}^{n} \left((A_0 + it)^{-1} + (A_0 - it)^{-1} \right) dt,
$$

According to [\[6,](#page-29-8) Proof of Thm. 4.2] for $f \in L^2(\mathbb{R})$ we have

$$
[J_0 f, f] = ||f||_2^2 + \frac{1}{\pi} \cdot \lim_{n \to \infty} \text{Re} \int_{1/n}^n \sqrt{\frac{2}{t}} \cdot [\phi_t, f][\phi_t, \overline{f}] dt, \tag{6.4}
$$

where $\phi_t(x) = e^{(i - \text{sgn}(x))x \cdot \sqrt{t/2}}$. Setting $f_1(x) = f(x)$ and $f_2(x) := f(-x)$, we obtain

$$
[\phi_t, f] = \int_0^\infty \left(e^{-\sqrt{\frac{t}{2}}(1-i)x} \overline{f_1(x)} - e^{-\sqrt{\frac{t}{2}}(1+i)x} \overline{f_2(x)} \right) dx,
$$

and thus

$$
[\phi_t, f][\phi_t, \overline{f}] = \sum_{j,k=1}^2 (-1)^{j+k} \int_0^{\infty} \int_0^{\infty} e^{-z_{j,k}(x,y)} \sqrt{\frac{t}{2}} \cdot \overline{f_j(x)} f_k(y) dx dy,
$$

where $z_{ik}(x, y) = x + y + i((-1)^{j}x + (-1)^{k}y)$. Taking into account that

$$
\lim_{n \to \infty} \int_{1/n}^{n} \sqrt{\frac{2}{t}} \cdot e^{-z\sqrt{t/2}} dt = \frac{4}{z}
$$

for any $z \in \mathbb{C}$ with Re $z > 0$, an application of Fubini's theorem and the dominated convergence theorem yields

$$
\lim_{n \to \infty} \int_{1/n}^{n} \sqrt{\frac{2}{t}} \cdot [\phi_t, f][\phi_t, \overline{f}] dt = 4 \sum_{j,k=1}^{2} (-1)^{j+k} \int_{0}^{\infty} \int_{0}^{\infty} \frac{\overline{f_j(x)} f_k(y)}{z_{j,k}(x, y)} dx dy.
$$

Taking real parts, we find that (see (6.4))

$$
[J_0 f, f] = ||f||_2^2 + \frac{2}{\pi} \iiint_0^{\infty} \left[\frac{\overline{f_1(x)} f_1(y) + \overline{f_2(x)} f_2(y)}{x + y} - \frac{2(x + y)}{x^2 + y^2} \operatorname{Re}(\overline{f_1(x)} f_2(y)) \right] dx dy.
$$

For *g*, $h \in L^2(\mathbb{R}^+)$ and a measurable symmetric kernel $L : \mathbb{R}^+ \times \mathbb{R}^+ \to [0, \infty)$ satisfying the homogeneity condition $L(tx, ty) = t^{-1}L(x, y)$ formally define the sesquilinear form

$$
\langle g, h \rangle_L := \int_0^\infty \int_0^\infty L(x, y) g(x) \overline{h(y)} \, dx \, dy.
$$

By [\[20](#page-29-27), Thm. 319], the form $\langle \cdot, \cdot \rangle_L$: $L^2(\mathbb{R}^+) \times L^2(\mathbb{R}^+) \to \mathbb{C}$ is well-defined and bounded if $||L||_* := \int_0^\infty L(x, 1)x^{-1/2} dx < \infty$. In this case, its bound is given by $||L||_*$. Here, we shall consider the kernels

$$
L_1(x, y) = \frac{1}{x + y}
$$
, $L_2(x, y) = \frac{2(x + y)}{x^2 + y^2}$, and $L_3 = 2L_1 + L_2$.

For these we have the bounds $||L_1||_* = \pi$, $||L_2||_* = 2\sqrt{2}\pi$, and $||L_3||_* = 2(1+\sqrt{2})\pi$. Thus,

$$
\begin{aligned} [J_0 f, f] &= \|f\|_2^2 + \frac{2}{\pi} (\langle f_1, f_1 \rangle_{L_1} + \langle f_2, f_2 \rangle_{L_1} - \text{Re} \, \langle f_1, f_2 \rangle_{L_2}) \\ &\le \|f\|_2^2 + 2 \big(\|f_1\|_{L^2(\mathbb{R}^+)}^2 + \|f_2\|_{L^2(\mathbb{R}^+)}^2 + 2\sqrt{2} \|f_1\|_{L^2(\mathbb{R}^+)} \|f_2\|_{L^2(\mathbb{R}^+)} \big) \\ &\le (3 + 2\sqrt{2}) \|f\|_2^2, \end{aligned}
$$

and therefore

$$
\tau_0 = ||JJ_0|| = \sup\{(JJ_0f, f)_2 : ||f||_2 = 1\} = \sup\{[J_0f, f] : ||f||_2 = 1\} \le 3 + 2\sqrt{2}.
$$

On the other hand, if we choose functions $f \in L^2(\mathbb{R})$ with $f_2 = -f_1$, we obtain

$$
[J_0 f, f] = ||f||_2^2 + \frac{2}{\pi} \int_0^\infty \int_0^\infty \left(\frac{2}{x+y} + \frac{2(x+y)}{x^2+y^2}\right) \overline{f_1(x)} f_1(y) dx dy
$$

= $||f||_2^2 + \frac{2}{\pi} \langle f_1, f_1 \rangle_{L_3}.$

For $\varepsilon > 0$, choosing a function $f_1 \in L^2(\mathbb{R}^+)$, $f_1 \ge 0$, with $||f_1||^2_{L^2(\mathbb{R}^+)} = \frac{1}{2}$ (i.e., $||f||_2 = 1$) and $\langle f_1, f_1 \rangle_{L_3} \ge \frac{1}{2} (||L_3||_* - \pi \varepsilon)$ leads to $[J_0 f, f] \ge 3 + 2\sqrt{2} - \varepsilon$, showing that

$$
\tau_0 = 3 + 2\sqrt{2}.
$$

Step 2. (*Estimation of the exceptional region K*) Let the operator *V* be defined by

$$
Vf := \operatorname{sgn} \cdot q \cdot f, \qquad \text{dom } V = \{ f \in L^2(\mathbb{R}) : qf \in L^2(\mathbb{R}) \}.
$$

By Lemma [6.1](#page-22-0) we have dom $A_0 \subset$ dom *V* and, for any $r > 0$ and $f \in$ dom A_0 ,

$$
||Vf||_2 \le (2r)^{1/p} \left(||f||_2 + \frac{1}{2\sqrt{3}\pi^2 pr^2} ||A_0f||_2 \right) ||q||_p.
$$

This implies that for all *r*, $\kappa > 0$ and all $f \in \text{dom } A_0$ we have

$$
||Vf||_2^2 \le ||q||_p^2 4^{\frac{1}{p}} (1 + \frac{1}{\kappa}) r^{\frac{2}{p}} ||f||_2^2 + \frac{(1+\kappa) 4^{\frac{1}{p}} ||q||_p^2}{12 \pi^4 p^2} r^{\frac{2}{p} - 4} ||A_0 f||_2^2. \tag{6.5}
$$

1

In particular, $A_0 + V$ is *J*-self-adjoint by the Kato-Rellich theorem. By [\[7,](#page-29-11) Thm. 1.1] the same is true for the operator *A*. And since dom $(A_0 + V) = H^2(\mathbb{R}) \subset D_{\text{max}} =$ dom *A*, it follows that $A_0 + V \subset A$ and thus $A_0 + V = A$.

From [\(6.5\)](#page-25-0) we get that with $\tau := \tau_0 = 3 + 2\sqrt{2}$ we have

$$
(1+\tau)\tau \, \|Vf\|_2^2 \le 2\alpha_{r,\kappa} \|f\|_2^2 + \beta_{r,\kappa} \|A_0f\|_2^2, \qquad f \in \text{dom}\, A_0,
$$

$$
\alpha_{r,\kappa} = \frac{(1+\tau)\tau}{2} \|q\|_p^2 4^{\frac{1}{p}} (1+\frac{1}{\kappa}) r^{\frac{2}{p}} \quad \text{and} \quad \beta_{r,\kappa} = \frac{(1+\tau)\tau (1+\kappa) 4^{\frac{1}{p}} \|q\|_p^2}{12 \pi^4 p^2} r^{\frac{2}{p}-4}.
$$

We have $\beta_{r,k} < \frac{\tau-1}{2\tau}$ if and only if $r > r_k := (\frac{\tau^2(1+\tau)(1+\kappa)4^{1/p}}{6(\tau-1)\pi^4 p^2} ||q||_p^2)^{1/(4-2/p)}$. Therefore, for $s > 1$ we set $r_{\kappa,s} := \left(\frac{\tau^2 (1+\tau)(1+\kappa) 4^{1/p}}{6(\tau-1)\pi^4 p^2} \|q\|_p^2 s \right)^{1/(4-2/p)}$ as well as

$$
a_{\kappa,s} := \alpha_{r_{\kappa,s},\kappa} = \frac{(1+\tau)\tau}{2} \left(\frac{8\tau^2(1+\tau)}{3(\tau-1)\pi^4 p^2} \right)^{\frac{1}{2p-1}} \|q\|_p^{\frac{4p}{2p-1}} (1+\frac{1}{\kappa})(1+\kappa)^{\frac{1}{2p-1}} s^{\frac{1}{2p-1}},
$$

$$
b_s := \beta_{r_{\kappa,s},\kappa} = \frac{\tau-1}{2\tau \cdot s} \quad \text{and} \quad \gamma_{\kappa,s} := \sqrt{\frac{1+\tau}{2\tau}} a_{\kappa,s}.
$$

The minimum of $(1 + \frac{1}{\kappa})(1 + \kappa)^{\frac{1}{2p-1}}$ is attained at $\kappa = 2p - 1$ and equals $\frac{2p}{2p-1}(2p)^{1/(2p-1)}$. We choose this *κ* and obtain $a_s = M_ps^{1/(2p-1)}$, where

$$
M_p := \frac{(1+\tau)\tau}{2} \left(\frac{16\tau^2(1+\tau)}{3(\tau-1)\pi^4 p}\right)^{\frac{1}{2p-1}} \frac{2p}{2p-1} \|q\|_p^{\frac{4p}{2p-1}} = \frac{4\tau^2}{1+\tau} s_p^{-\frac{1}{2p-1}} \|q\|_p^{\frac{4p}{2p-1}} C_p^2,
$$

as well as $\gamma_s = \sqrt{\frac{1+\tau}{2\tau}} a_s$. Now, the second part of Theorem [5.2](#page-18-1) implies that for each *s* > 1 the operator $A = A_0 + V$ is *J*-non-negative over $\overline{C} \setminus K_s$, where

$$
K_s := \bigcup_{t \in [-\gamma_s, \gamma_s]} B_{\sqrt{\frac{1+\tau}{2\tau(1-b_s)}(a_s+b_s t^2)}}(t).
$$

For any $\lambda \in K_s$ we have

$$
|\operatorname{Im} \lambda|^2 \leq \frac{1+\tau}{2\tau(1-b_s)}(a_s+b_s\gamma_s^2) = \frac{1+\tau}{2\tau-\frac{\tau-1}{s}}\left(1+\frac{\tau^2-1}{4\tau^2s}\right)a_s = \frac{1+\tau}{4\tau^2}M_p f(s)^2 s^{\frac{1}{2p-1}}.
$$

The minimum of $s \mapsto f(s)^2 s^{1/(2p-1)}$ on $(1, \infty)$ is attained at s_p . Choosing $s = s_p$, we find that *A* is *J*-non-negative over $\overline{\mathbb{C}} \setminus K'$, where $K' := K_{s_p}$, and we have just proved the claimed bound on $|\text{Im }\lambda|$ for $\lambda \in K'$. Furthermore, for $\lambda \in K'$ we have

$$
|\operatorname{Re}\lambda| \leq \gamma_{s_p} + \sqrt{\frac{(1+\tau)(a_{s_p} + b_{s_p}\gamma_{s_p}^2)}{2\tau(1-b_{s_p})}} = \sqrt{\frac{1+\tau}{2\tau}} \sqrt{\frac{4\tau^2}{1+\tau}} \|q\|_p^{\frac{4p}{2p-1}} C_p^2 + C_p f(s_p) \|q\|_p^{\frac{2p}{2p-1}}
$$

= $C_p(\sqrt{2\tau} + f(s_p)) \|q\|_p^{\frac{2p}{2p-1}},$

and the theorem is proved.

Remark 6.3 (a) The bound on the real part in (6.3) can be further slightly improved by minimizing the expression $\gamma_s + \sqrt{\frac{1+\tau}{2\tau(1-b_s)}(a_s + b_s\gamma_s^2)}$, where $\tau = 3 + 2\sqrt{2}$.

Fig. 5 Comparison of the enclosures in (6.3) (blue) and (6.6) (red). The left figure depicts the bounds on the imaginary part in dependence of p, respectively, whereas the right figure compares the bound on | λ | in (6.6 Fig. 5 Comparison of the enclosures in [\(6.3\)](#page-23-2) (blue) and [\(6.6\)](#page-28-0) (red). The left figure depicts the bounds on the imaginary part in dependence of p, respectively, whereas the right figure compares the bound on $|\lambda|$ in (6.6)

(b) In [\[6\]](#page-29-8) it was proved that $\tau_0 \leq 9$. In the proof of Theorem [6.2](#page-23-0) we have now shown that $\tau_0 = 3 + 2\sqrt{2}$.

(c) Estimates on the non-real spectrum of singular indefinite Sturm-Liouville operators have been obtained in [\[7\]](#page-29-11) for various weights and potentials. In the case of the signum function as weight and a negative potential $q \in L^p(\mathbb{R})$ the enclosure in [\[7,](#page-29-11) Cor. 2.7 (ii)] for the non-real eigenvalues λ of A reads as follows:

$$
|\operatorname{Im}\lambda| \le 2^{\frac{2p+1}{2p-1}} \cdot 3 \cdot \sqrt{3} \|q\|_p^{\frac{2p}{2p-1}} \quad \text{and} \quad |\lambda| \le \left(2^{\frac{2p+1}{2p-1}} \cdot 3 \cdot \sqrt{3} + 2^{\frac{3-2p}{2p-1}} \cdot 9\right) \|q\|_p^{\frac{2p}{2p-1}}.
$$
\n
$$
(6.6)
$$

A direct comparison shows that the enclosure for the non-real spectrum of *A* in The-orem [6.2](#page-23-0) is strictly better in the sense that the region K in (6.3) is properly contained in that described by [\(6.6\)](#page-28-0) (see Fig. [5\)](#page-27-0). This is remarkable inasmuch as our bound [\(6.3\)](#page-23-2) was mainly obtained by applying the abstract perturbation result Theorem [5.2,](#page-18-1) whereas in [\[7\]](#page-29-11) the authors work directly with the differential expressions. It is also noteworthy that the estimates in [\[7,](#page-29-11) Cor. 2.7 (ii)] are of the same form $C||q||_p^{2p/(2p-1)}$ as in [\(6.3\)](#page-23-2).

(d) Recently, the bound on the imaginary part of the eigenvalues of *A* could be further significantly improved in [\[9](#page-29-12)] by using a Birman-Schwinger type principle. However, the bounding region in [\[9\]](#page-29-12) is not compact. To be precise, it was shown that each eigenvalue λ of *A* in [\(6.1\)](#page-21-1) satisfies $2^{\frac{3}{2p}-1} |\lambda|^{\frac{1}{p}} |\operatorname{Im} \lambda|^{1-\frac{1}{p}} \leq (|\lambda| +$ $|\text{Re }\lambda|\big)^{\frac{1}{2p}}\|q\|_p$, which implies

$$
|\operatorname{Im} \lambda| \leq 2 \left(\frac{3\sqrt{3}}{16} \right)^{\frac{1}{2p-1}} \|q\|_p^{\frac{2p}{2p-1}}.
$$

Acknowledgements The author thanks Jussi Behrndt, Christian Gérard, David Krejčiřík, Ilya Krishtal, Christiane Tretter, and Carsten Trunk for valuable discussions and hints.

Funding Open Access funding enabled and organized by Projekt DEAL.

Declarations

Data Availability Data sharing is not applicable to this article as no datasets were generated or analyzed during the current study.

Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit [http://creativecommons.org/licenses/by/4.0/.](http://creativecommons.org/licenses/by/4.0/)

References

- 1. Azizov, T.Ya., Behrndt, J., Philipp, F., Trunk, C.: On domains of powers of linear operators and finite rank perturbations. Oper. Theory: Adv. Appl. **188**, 31–36 (2008)
- 2. Albeverio, S., Motovilov, A.K., Shkalikov, A.A.: Bounds on variation of spectral subspaces under J-self-adjoint perturbations. Integral Equations Oper. Theory **64**, 455–486 (2009)
- 3. Albeverio, S., Motovilov, A.K.: and Christiane Tretter, Bounds on the spectrum and reducing subspaces of a *J* -self-adjoint operator. Indiana Univ. Math. J. **59**, 1737–1776 (2010)
- 4. Behrndt, J., Jonas, P.: On compact perturbations of locally definitizable self-adjoint relations in Krein spaces. Integral Equations Oper. Theory **52**, 17–44 (2005)
- 5. Behrndt, J., Philipp, F.: Spectral analysis of singular ordinary differential operators with indefinite weights. J. Differ. Equ. **248**, 2015–2037 (2010)
- 6. Behrndt, J., Philipp, F., Trunk, C.: Bounds on the non-real spectrum of differential operators with indefinite weights. Math. Ann. **357**, 185–213 (2013)
- 7. Behrndt, J., Schmitz, P., Trunk, C.: Spectral bounds for indefinite singular Sturm-Liouville operators with uniformly locally integrable potentials. J. Differ. Equ. **267**, 468–493 (2019)
- 8. Bender, C.M., Boettcher, S.: Real spectra in non-Hermitian Hamiltonians having PT-symmetry. Phys. Rev. Lett. **80**, 5243–5246 (1998)
- 9. Cuenin, J.-C., Ibrogimov, O.O.: Sharp spectral bounds for complex perturbations of the indefinite Laplacian. J. Funct. Anal. **280**, 108804 (2021)
- 10. Cuenin, J.-C., Tretter, C.: Non-symmetric perturbations of self-adjoint operators. J. Math. Anal. Appl. **441**, 235–258 (2016)
- 11. Curgus, B.: On the regularity of the critical point infinity of definitizable operators. Integral Equations ´ Oper. Theory **8**, 462–488 (1985)
- 12. Curgus, B., Langer, H.: A Krein space approach to symmetric ordinary differential operators with an ´ indefinite weight function. J. Differ. Equ. **79**, 31–61 (1989)
- 13. Curgus, B., Najman, B.: The operator $sgn(x) \frac{d^2}{dx^2}$ is similar to a selfadjoint operator in $L^2(\mathbb{R})$. Proc. Am. Math. Soc. **123**, 1125–1128 (1995)
- 14. Eckhardt, J., Kostenko, A.: An isospectral problem for global conservative multi-peakon solutions of the Camassa-Holm equation. Comm. Math. Phys. **329**, 893–918 (2014)
- 15. Edmunds, D.E., Evans, W.D.: Spectral theory and differential operators, Oxford Mathematical Monographs. The Clarendon Press Oxford University Press, Oxford Science Publications, New York (1987)
- 16. Folland, G.B.: Real Analysis. John Wiley & Sons, Inc., New York, Chichester, Brisbane, Toronto, Singapore (1984)
- 17. Gérard, C.: Scattering theory for Klein-Gordon equations with non-positive energy. Ann. Henri Poincaré **13**, 883–941 (2012)
- 18. Georgescu, V., Gérard, C., Häfner, D.: Resolvent and propagation estimates for Klein-Gordon equations with non-positive energy. J. Spectr. Theory **5**, 113–192 (2015)
- 19. Giribet, J., Langer, M., Martínez-Pería, F., Philipp, F., Trunk, C.: Spectral enclosures for a class of block operator matrices. J. Funct. Anal. **278**(10), 108455 (2020)
- 20. Hardy, G.H., Littlewood, J.E., Pólya, G.: Inequalities. Cambridge University Press, Cambridge (1934)
- 21. Jacob, B., Tretter, C., Trunk, C., Vogt, H.: Systems with strong damping and their spectra. Math. Methods Appl. Sci. **41**, 6546–6573 (2018)
- 22. Jonas, P.: Über die Erhaltung der Stabilität *J* -positiver Operatoren bei *J* -positiven und *J* -negativen Störungen. Math. Nachr. **65**, 211–218 (1975)
- 23. Jonas, P.: On a class of self-adjoint operators in Krein space and their compact perturbations. Integral Equations Oper. Theory **11**, 351–384 (1988)
- 24. Jonas, P.: On the spectral theory of operators associated with perturbed Klein-Gordon and wave type equations. J. Oper. Theory **29**, 207–224 (1993)
- 25. Jonas, P.: Riggings and relatively form bounded perturbations of nonnegative operators in Krein spaces. Oper. Theory: Adv. Appl. **106**, 259–273 (1998)
- 26. Jonas, P.: On locally definite operators in Krein spaces, in: Spectral Theory and its Applications. Ion Colojoară Anniversary Volume, Theta 95-127 (2003)
- 27. Jonas, P., Langer, H.: Some questions in the perturbation theory of *J* -nonnegative operators in Krein spaces. Math. Nachr. **114**, 205–226 (1983)
- 28. Kato, T.: Perturbation theory for linear operators. Springer-Verlag, Berlin Heidelberg (1995)
- 29. Kong, Q., Möller, M., Wu, H., Zettl, A.: Indefinite Sturm-Liouville problems. Proceedings of the Royal Society of Edinburgh **133**(A), 639–652 (2003)
- 30. Krein, M.G., Langer, H.: On some mathematical principles in the linear theory of damped oscillations of continua I. Integral Equ. Oper. Theory **1**, 364–399 (1978)
- 31. Krein, M.G., Langer, H.: On some mathematical principles in the linear theory of damped oscillations of continua II. Integral Equ. Oper. Theory **1**, 539–566 (1978)
- 32. Krejčiřík, D., Siegl, P.: Elements of spectral theory without the spectral theorem. In: Bagarello, F., Gazeau, J.P., Szafraniec, F.H., Znojil, M. (eds.) Non-selfadjoint operators in quantum physics: Mathematical aspects. John Wiley & Sons, Inc., Hoboken, New Jersey (2015)
- 33. Lancaster, P., Markus, A.S., Matsaev, V.I.: Definitizable operators and quasihyperbolic operator polynomials. J. Funct. Anal. **131**, 1–28 (1995)
- 34. Langer, H.: Spectral functions of definitizable operators in Krein spaces. Lect. Notes Math. **948**, 1–46 (1982)
- 35. Langer, H., Langer, M., Markus, A., Tretter, C.: The Virozub-Matsaev condition and spectrum of definite type for self-adjoint operator functions. Compl. anal. oper. theory **2**, 99–134 (2008)
- 36. Langer, H., Markus, A., Matsaev, V.: Locally definite operators in indefinite inner product spaces. Math. Ann. **308**, 405–424 (1997)
- 37. Langer, H., Markus, A., Matsaev, V.: Linearization and compact perturbation of self-adjoint analytic operator functions. Oper. Theory.: Adv. Appl. **118**, 255–285 (2000)
- 38. Langer, H., Markus, A., Matsaev, V.: Self-adjoint analytic operator functions and their local spectral function. J. Funct. Anal. **235**, 193–225 (2006)
- 39. Langer, H., Markus, A., Matsaev, V.: Self-adjoint analytic operator functions: local spectral function and inner linearization. Integr. equ. oper. theory **63**, 533–545 (2009)
- 40. Langer, H., Najman, B., Tretter, C.: Spectral theory of the Klein-Gordon equation in Pontryagin spaces. Comm. Math. Phys. **267**, 159–180 (2006)
- 41. Langer, H., Najman, B., Tretter, C.: Spectral theory of the Klein-Gordon equation in Krein spaces. Proc. Edinb. Math. Soc. **51**, 1–40 (2008)
- 42. Langer, H., Tretter, C.: A Krein space approach to PT-symmetry. Czech. J. Phys. **54**, 1113–1120 (2004)
- 43. Langer, H., Tretter, C.: Corrigendum to: A Krein space approach to PT-symmetry. Czech. J. Phys. **56**, 1063–1064 (2006)
- 44. Philipp, F.: Locally definite normal operators in Krein spaces. J. Funct. Anal. **262**, 4929–4947 (2012)
- 45. Philipp, F.: Indefinite Sturm-Liouville operators with periodic coefficients. Oper. Matrices **7**, 777–811 (2013)
- 46. Tanaka, T.: PT-symmetric quantum theory defined in a Krein space. J. Phys. A: Math. Gen. **39**, 369–376 (2006)
- 47. Tretter, C.: Spectral Theory of Block Operator Matrices and Applications. Imperial College Press, London (2008)
- 48. Zettl, A.: Sturm-Liouville Theory. AMS, Providence (2005)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.